

Smoothed Particle Hydrodynamics, fact checking: from theory to applications

David LE TOUZÉ

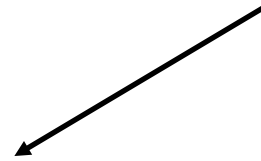
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For a start: SPH in 3 words



Smoothed Particle Hydrodynamics

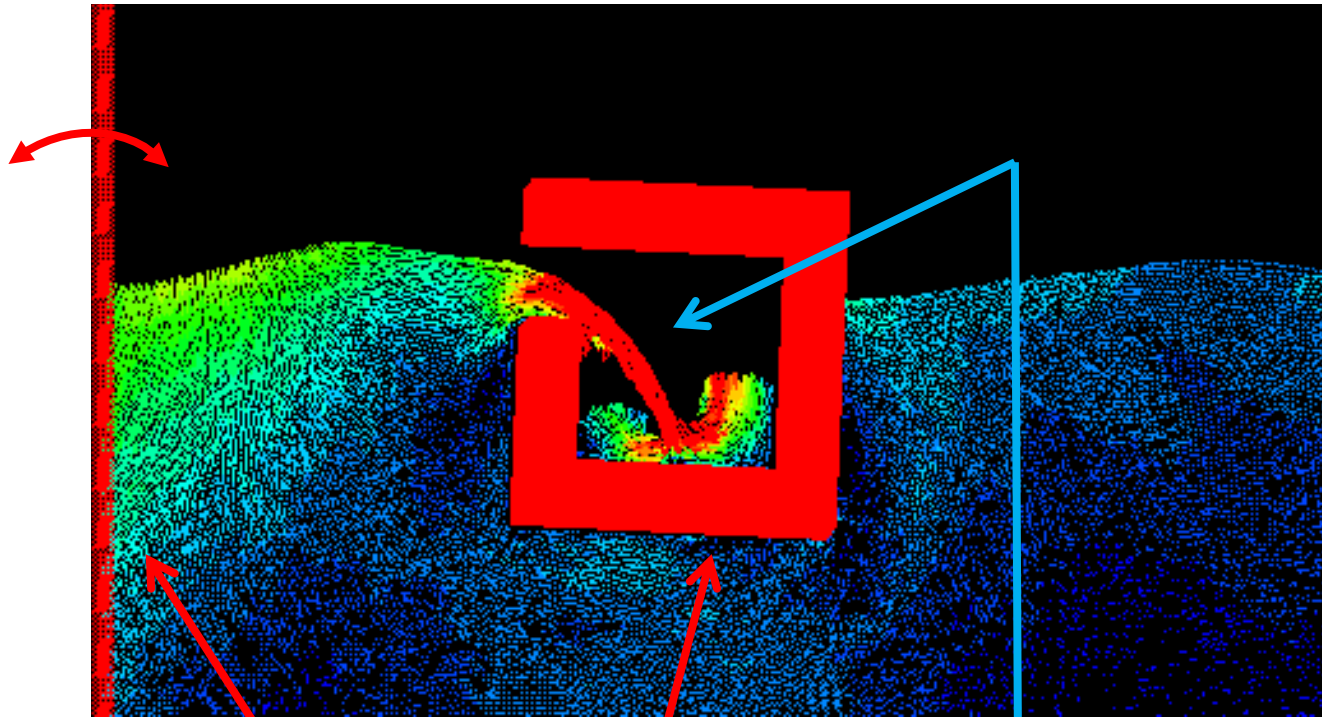


A computational method for solving continuum mechanics problems...

... with large deformations, multiple objects and complex interfaces

Large deformations and complex interfaces...

What we were doing in 2000



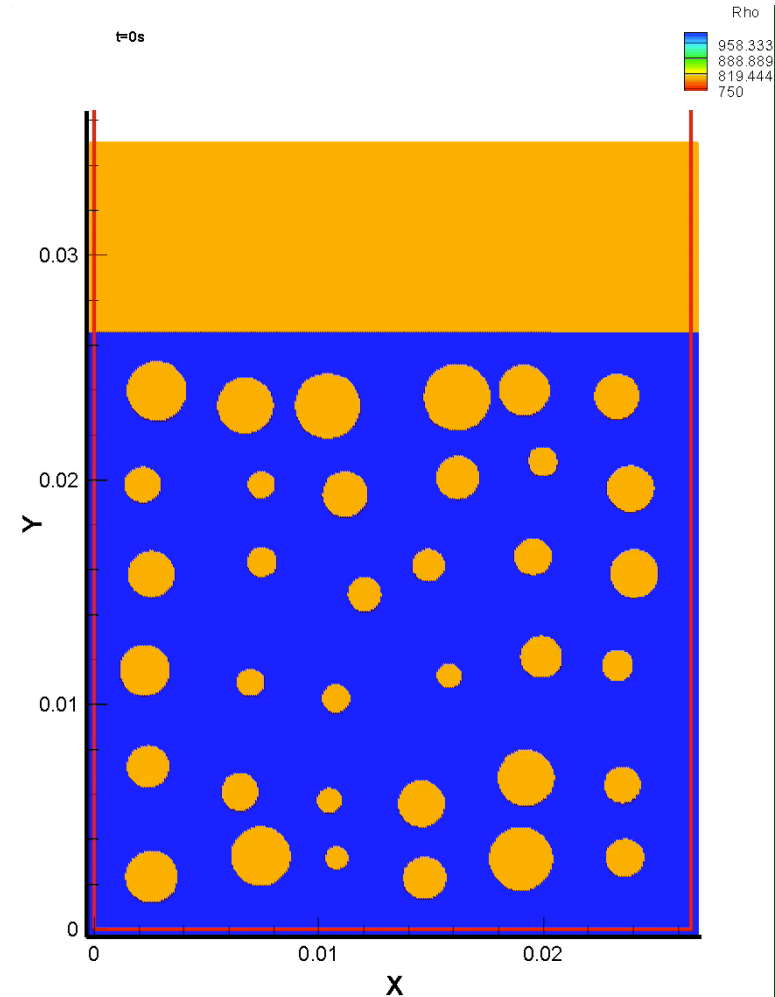
2 moving bodies = wavemaker + floating body + 1 complex interface

Large deformations and complex interfaces...

What we were doing 10 years ago



Tire aquaplaning



Water-oil separation

Large deformations and complex interfaces...

What we were doing 5 years ago

HYDROCEAN

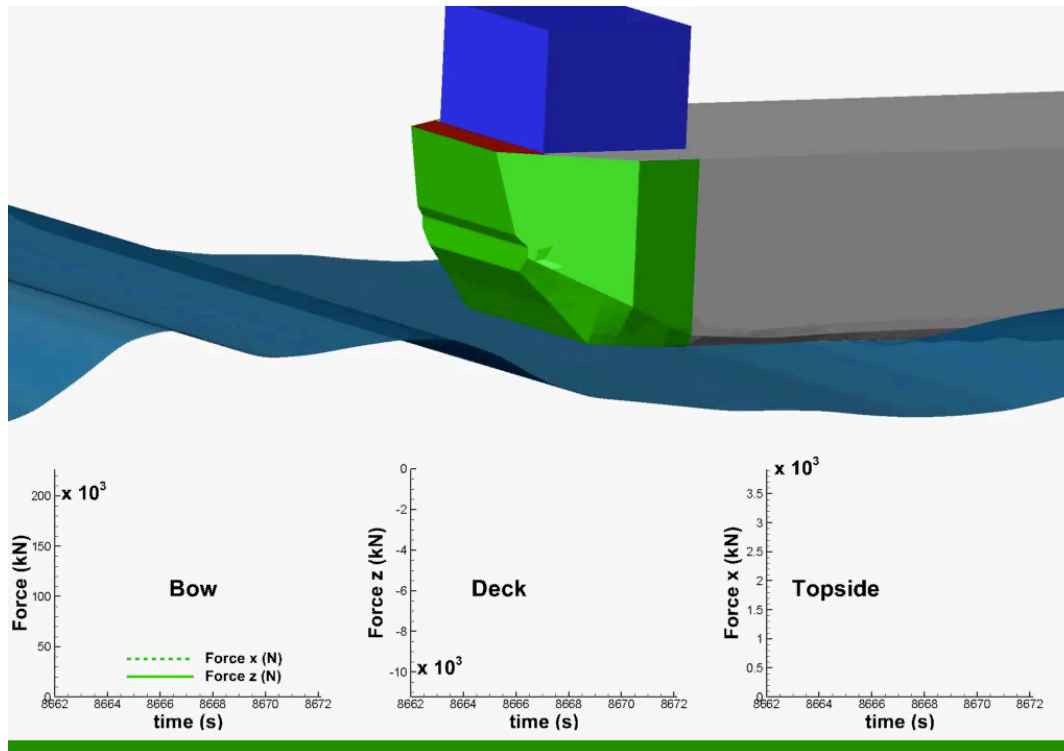
SPH-flow



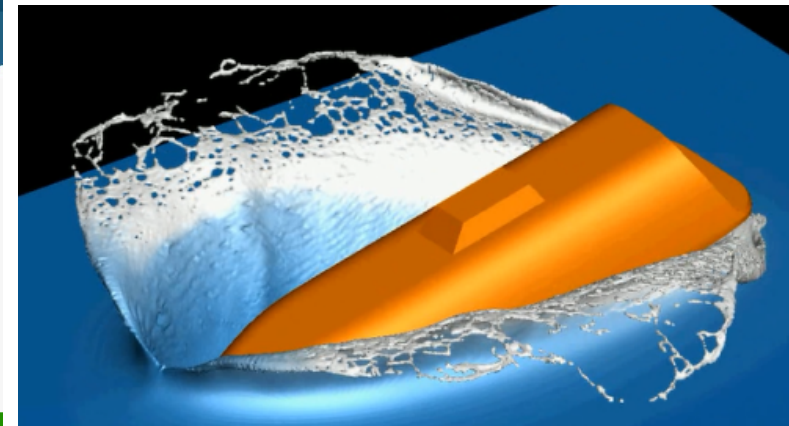
TOTAL
Technip

ECN
Centrale
Nantes

Focused wave on a FLNG
- 20 500 000 particles
- 2048 cpus



Greenwater loads on a ship deck

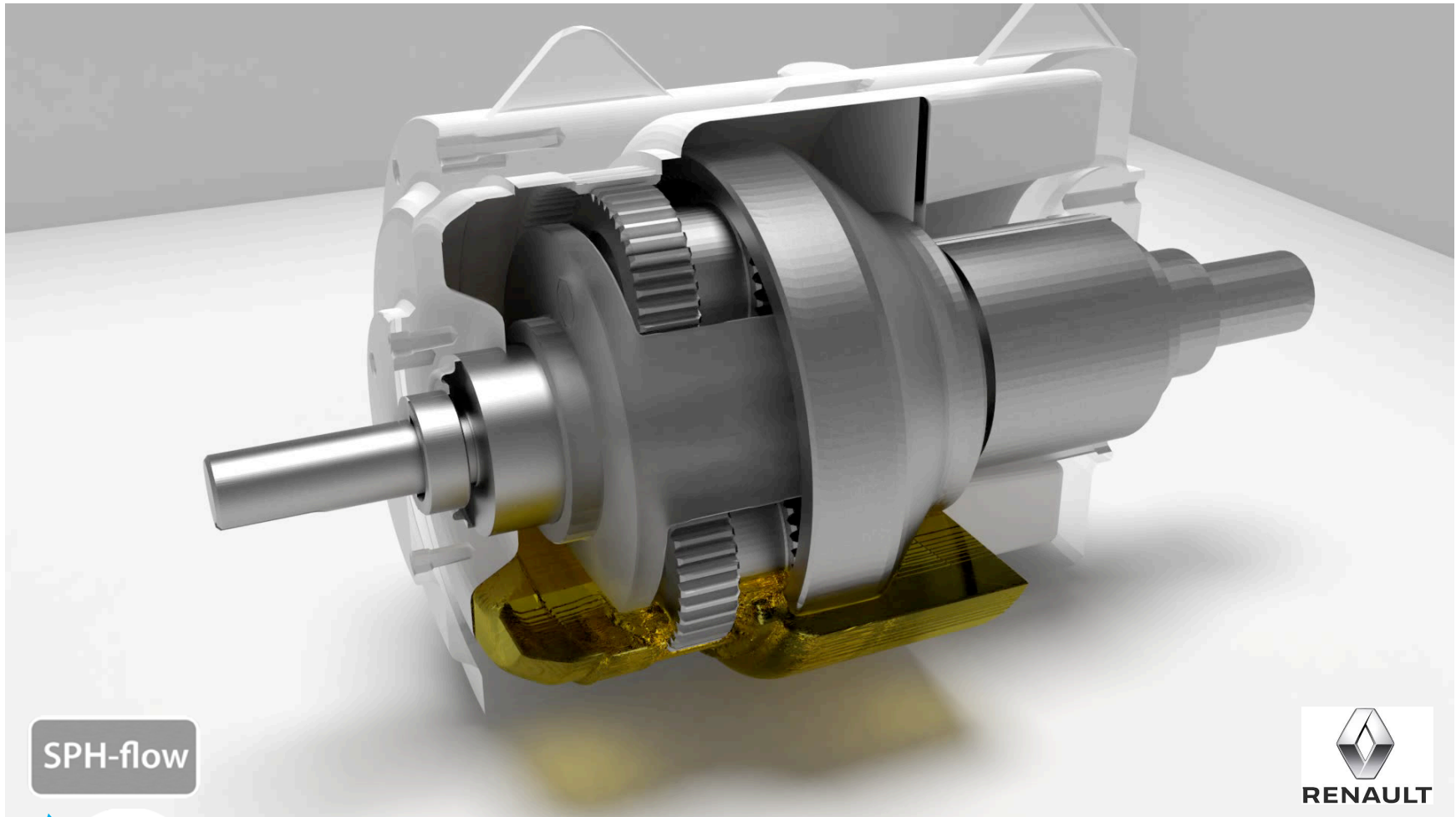


Lifeboat emergency ditching

Large deformations and complex interfaces...



What we are doing now



SPH-flow



Gearbox lubrication

For a start: SPH in 3 words

Smoothed Particle Hydrodynamics



Complex interfaces = large deformations, fragmentation, coalescence...

+

Complex (multiple body) motions = small gaps between objects, contact...

=

Very difficult for mesh-based methods, especially if we want accurate results

⇒ follow the interfaces/motions = Lagrangian

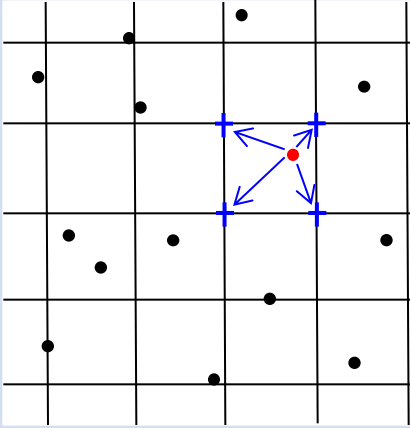
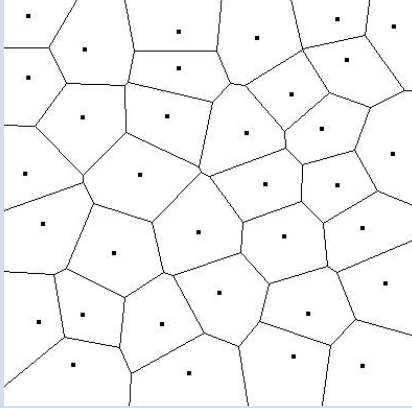
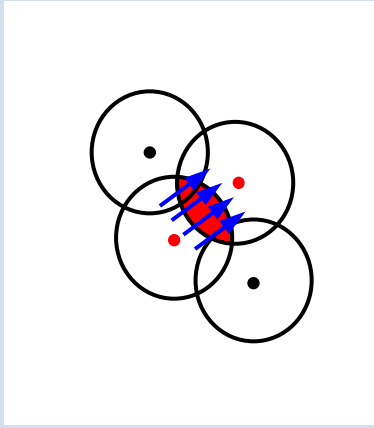
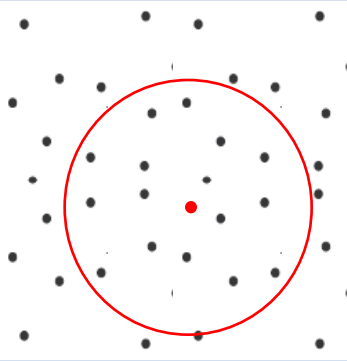
+

⇒ not possible to use a mesh = meshless

=

particle method

Meshless?

Projection	Tesselation	Face construction	Truly meshless
			
Particle-Mesh ...	Voronoi-FVM Particle-FEM ...	FVPM ...	SPH MPS ...

Meshless?



Truly meshless

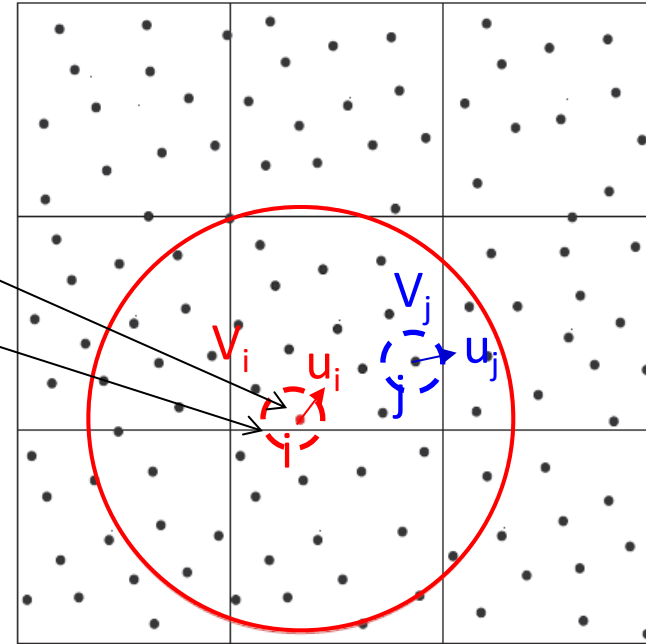
- + : any configuration can be easily described
- : no description of how the volume of a particle is spread around its location, and on how it will deform in time

Partly meshless (projection/reconstruction)

- + : easier to define convergent operators / make mathematical analysis of the schemes
- : less general / complex and costly implementation / **how to treat interfaces?**
especially free-surface?

Particle

Particle: a volumic element (volume V_i) of barycenter the scattered point i



Particle method:

=

meshless (i.e. no connectivity, NOT no space discretization)

+

Lagrangian (material evolution : $d\cdot/dt = \dots$): particle i evolves at its material speed \mathbf{u}_i

For a start: SPH in 3 slides

Smoothed Particle Hydrodynamics



E.g., for Navier-Stokes, in Lagrangian (material) formalism

$$\frac{d\mathbf{x}}{dt} = \mathbf{u}$$

$$\frac{d\rho}{dt} = -\rho \operatorname{div} \mathbf{u}$$

$$\frac{d\mathbf{u}}{dt} = \mathbf{g} - \frac{\operatorname{grad} p}{\rho} + \nu \Delta \mathbf{u} + \frac{\nu}{3} \operatorname{grad}(\operatorname{div} \mathbf{u})$$

$$p = f(\rho)$$

=> how do we calculate the spatial operators with no mesh?

For a start: SPH in 3 slides

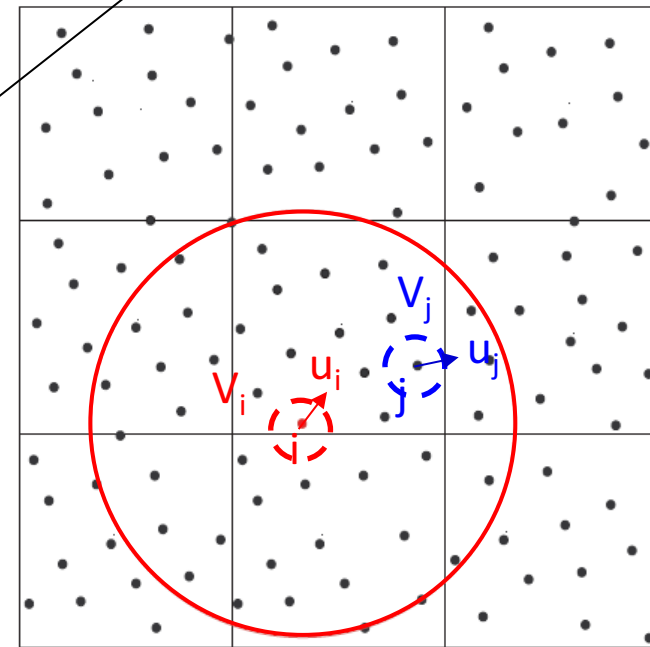
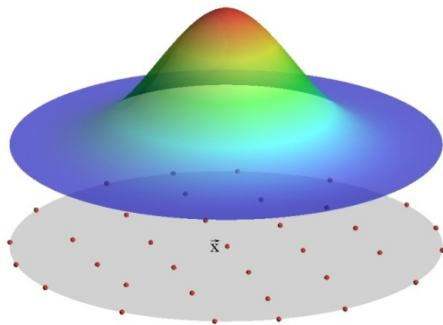
=> mollification + discretisation

$$\langle \mathbf{u} \rangle_i = \int \mathbf{u}(\mathbf{x}) \delta(|\mathbf{x} - \mathbf{x}_i|) d\mathbf{x} \stackrel{\textcircled{1}}{\cong} \int \mathbf{u}(\mathbf{x}) W(|\mathbf{x} - \mathbf{x}_i|) d\mathbf{x} \stackrel{\textcircled{2}}{\cong} \sum_j \mathbf{u}_j W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$

Known at the
neighbour
scattered points
(particle)

Volume associated to
the scattered points
(particles)

Analytical
function
(kernel)



Compact support!!!

+ transfer of the differentiation from the field
to the kernel to get differential operators:

$$\nabla f(\vec{r}) \stackrel{\textcircled{1}}{\approx} \int_D \nabla f(\vec{x}) W(\vec{r} - \vec{x}) d\vec{x} = \int_D f(\vec{x}) \nabla W(\vec{r} - \vec{x}) d\vec{x}$$

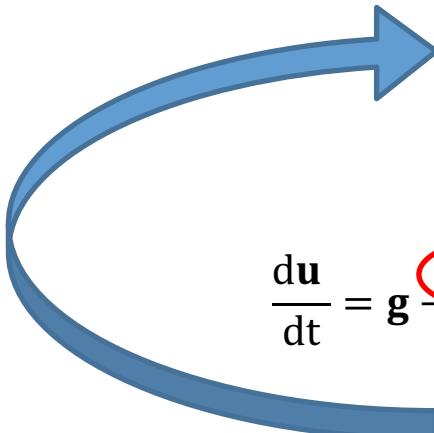
analytical

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Smoothed Particle Hydrodynamics



E.g., for Navier-Stokes, in Lagrangian (material) formalism



$$\frac{dx}{dt} = \mathbf{u}$$

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$$p = f(\rho)$$

Explicit time integration

For a start: SPH in 3 slides

Smoothed Particle Hydrodynamics

A computational method for **solving continuum mechanics** problems...

... with large deformations, multiple objects and complex interfaces

⇒ we are doing **engineering**, not movies or games

⇒ we want **accurate** stresses (e.g. **pressure**), **forces**, deformations...



SPH fact checking



partly true

1. SPH cannot be solved implicitly (as mesh-based) / it has very small time steps

1. Fully-implicit? Small time steps?

A **fully implicit** scheme would mean that time derivatives are expressed using the solution at next time step (n+1), including the displacements:

$$\frac{d\vec{x}}{dt} = \vec{V}^{(n+1)}$$

⇒ This would lead to a complex implementation in practice, and attempts we made showed that the resulting scheme is too diffusive

⇒ No one does that in SPH

⇒ 2 ways : **fully-explicit** OR, at best, **semi-implicit**

⇒ Fully explicit = WCSPH

$$\begin{aligned} \frac{dx}{dt} &= \mathbf{u} \\ \frac{d\rho}{dt} &= -\rho \operatorname{div} \mathbf{u} \\ \frac{d\mathbf{u}}{dt} &= \mathbf{g} - \frac{\operatorname{grad} p}{\rho} + \nu \Delta \mathbf{u} + \frac{\nu}{3} \operatorname{grad}(\operatorname{div} \mathbf{u}) \\ p &= f(\rho) \end{aligned}$$

⇒ **Stability criteria**: $\Delta t < \Delta x / c_s$

Semi-implicit = ISPH (or MPS)

$$\begin{aligned} \frac{dx}{dt} &= \mathbf{u} \\ \operatorname{div} \mathbf{u} &= 0 \\ \frac{d\mathbf{u}}{dt} &= \mathbf{g} - \frac{\operatorname{grad} p}{\rho} + \nu \Delta \mathbf{u} \end{aligned}$$

$\Delta t < \Delta x / |\mathbf{u}|_{\max}$

SPH fact checking



partly true

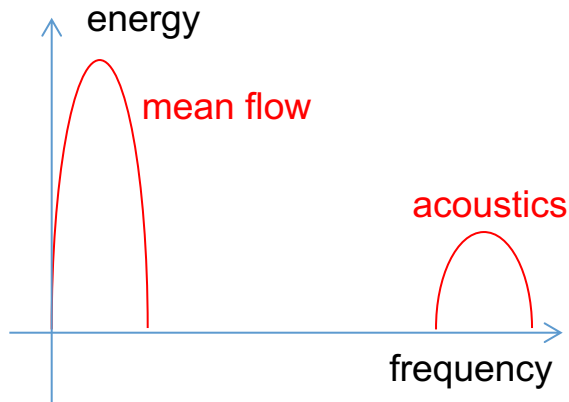
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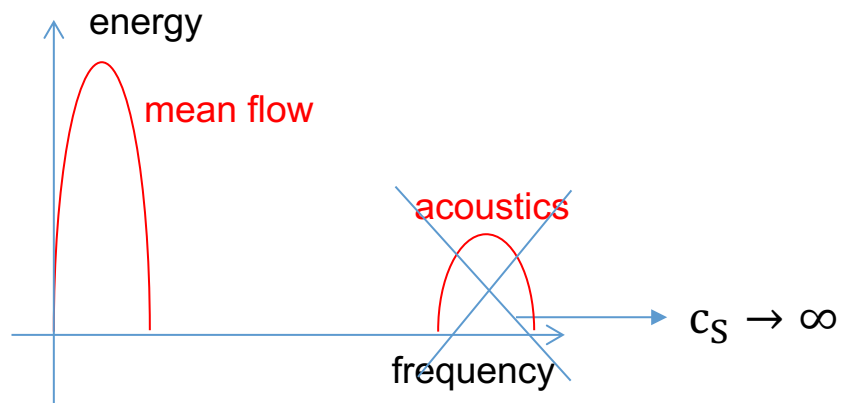
2. Weak-compressibility is an unphysical trick

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Low Mach physical situation = acoustics is superimposed to the incompressible part of the flow, and fully separated in frequencies

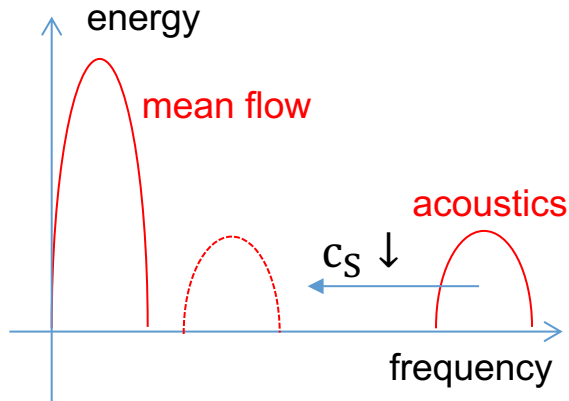


=> Incompressible assumption: all waves have infinite speed

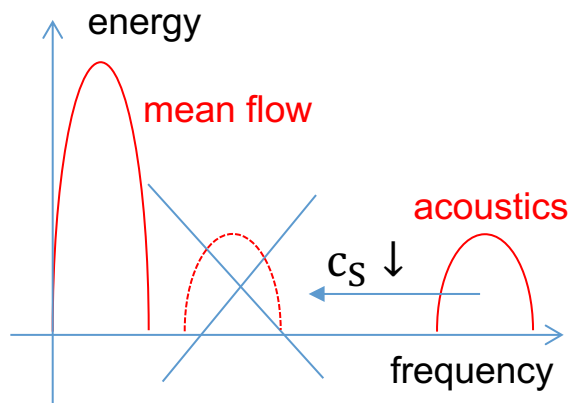


2. Weak-compressibility is an unphysical trick

=> **Weakly-compressible assumption**: we can change the sound speed since we are not interested in the acoustic part of the flow (at low Mach)



=> N.B. : **Weakly-compressible assumption + filtering**



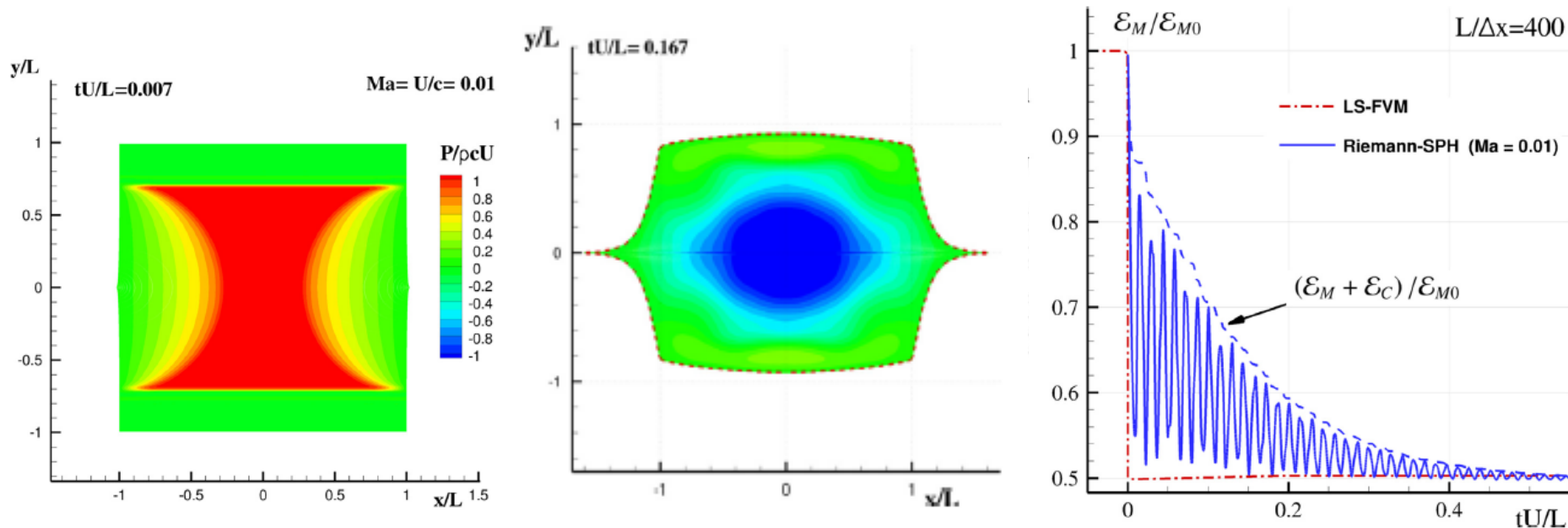
2. Weak-compressibility is an unphysical trick

⇒ **Weakly-compressible assumption**, summarizing:

- It is not a bigger assumption than supposing the flow incompressible
- Without filtering it implies that physical pressure oscillations should be present in the solution (at unphysical frequencies)
- **With perfect filtering, it is equivalent to the incompressible assumption**
- It permits to lower the sound speed provided Ma stays lower than 0.1 at least, i.e. $c_s > 10|\mathbf{u}|_{\max}$
- It thus permits to loosen the acoustic CFL stability condition: $\Delta t < \frac{\Delta x}{(10|\mathbf{u}|_{\max})}$
- This induces only a **factor 10 with ISPH**, $\Delta t < \Delta x/|\mathbf{u}|_{\max}$
(N.B. : ISPH requires imposing conditions at the free surface + solving a system at each time step)

2. Weak-compressibility is an unphysical trick

=> we checked that **weakly-compressible solution matches theoretical incompressible one** once acoustic oscillations are damped, **even at impact**



Le Touzé D. et al., A critical investigation of smoothed particle hydrodynamics applied to problems with free-surfaces, *Int. J. Numer. Meth. Fluids* **73**, 2013

Marrone S. et al., Prediction of energy losses in water impacts using incompressible and weakly compressible models, *J. Fluid Struct.* **54**, 2015

SPH fact checking



partly true

1. SPH cannot be solved implicitly (as mesh-based) / it has very small time steps

partly true

2. Weak-compressibility is an unphysical trick

not fully true

3. SPH conserves everything

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Let's restrict to Euler equations for now (perfect fluid)

$$\frac{d\mathbf{x}}{dt} = \mathbf{u}$$

$$\frac{d\rho}{dt} = -\rho \operatorname{div} \mathbf{u}$$

$$\frac{d\mathbf{u}}{dt} = \mathbf{g} - \frac{\operatorname{grad} p}{\rho}$$

$$p = f(\rho)$$

for **Momentum** conservation (action/reaction)

$$\langle \nabla p \rangle_i = \sum_j (p_j + p_i) W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$

$$\langle \operatorname{div} \mathbf{u} \rangle_i = \sum_j (\mathbf{u}_j - \mathbf{u}_i) \cdot \nabla W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$

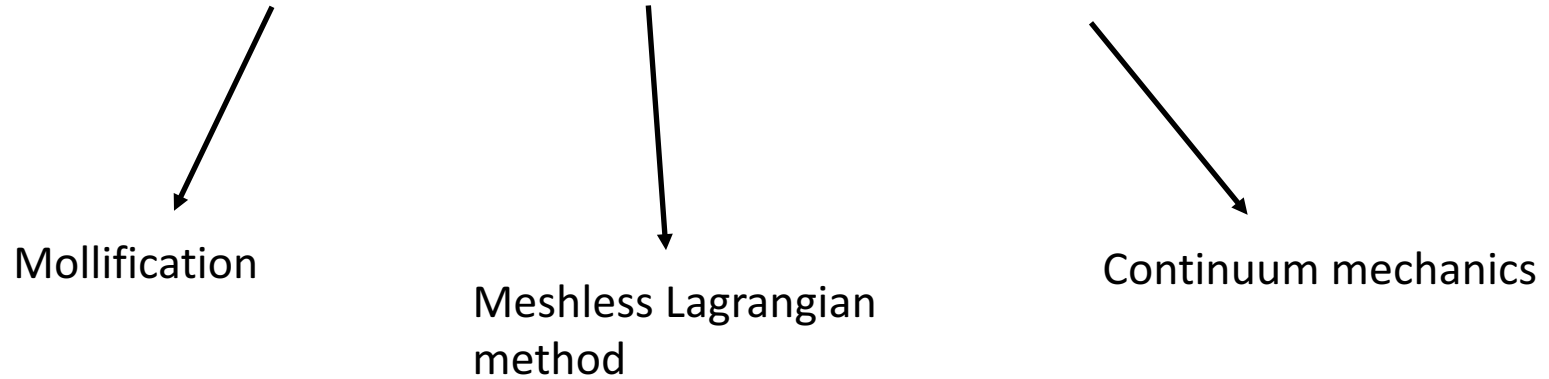
Mass is conserved (particle method)

for **Energy** conservation (Hamiltonian)

Total volume is not conserved!

3. SPH conserves everything

Smoothed Particle Hydrodynamics



Vision 1 (original) = Hamiltonian mechanics of a discrete system

A **fluid dynamics** problem numerically seen as a system of particles using a **mollification**.

AND

A **meshless numerical method** for discretizing Lagrangian PDEs.

Vision 2 (standard) = numerical methods for PDEs describing continuous media

SPH fact checking



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1. SPH cannot be solved implicitly (as mesh-based) / it has very small time steps

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no more true

4. SPH is unstable / SPH pressures are noisy

4. Stability? Pressure field quality?

Let's restrict to Euler equations for now (perfect fluid)

$$\frac{d\mathbf{x}}{dt} = \mathbf{u}$$

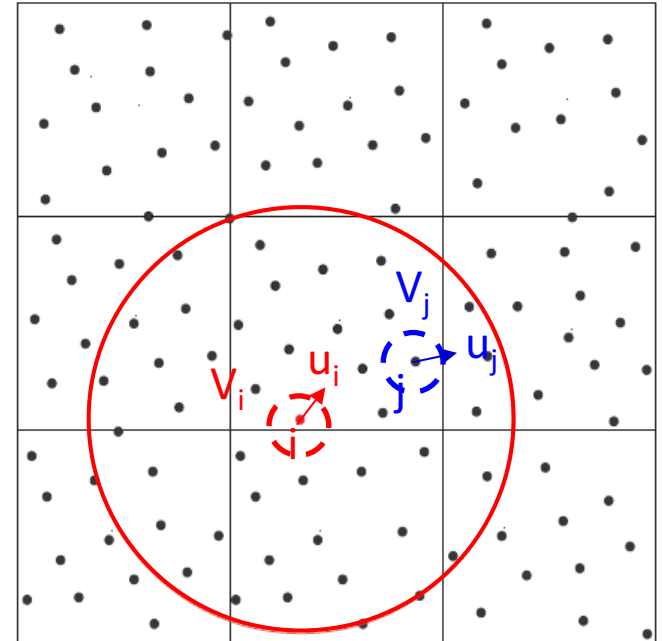
$$\frac{d\rho}{dt} = -\rho \operatorname{div} \mathbf{u}$$

$$\frac{d\mathbf{u}}{dt} = \mathbf{g} - \frac{\operatorname{grad} p}{\rho}$$

$$p = f(\rho)$$

$$\langle \nabla p \rangle_i = \sum_j (p_j + p_i) W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$

$$\langle \operatorname{div} \mathbf{u} \rangle_i = \sum_j (\mathbf{u}_j - \mathbf{u}_i) \cdot \nabla W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$



Scheme which is **centered in space** + **fully explicit** => need for stabilization

4. Stability? Pressure field quality?



centered + explicit => need for **stabilization**

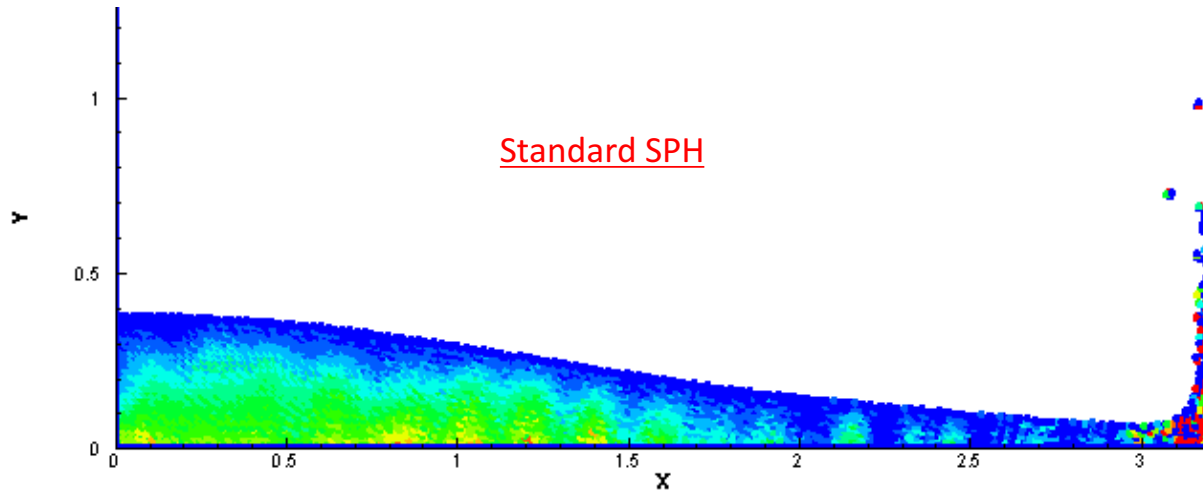
Method 1 (stable but not enough!)

last equations + **artificial viscosity** added through the pressure term (original Monaghan SPH scheme)

$$\rho_i \frac{d\vec{v}_i}{dt} = \rho_i \vec{g}_i - \sum_{j \in \Omega} (P_i + P_j) \vec{\nabla}_i W_{ij} \omega_j + \alpha h c_0 \rho_0 \sum_{j \in \Omega} \Pi_{ij} \vec{\nabla} W_{ij} \omega_j,$$

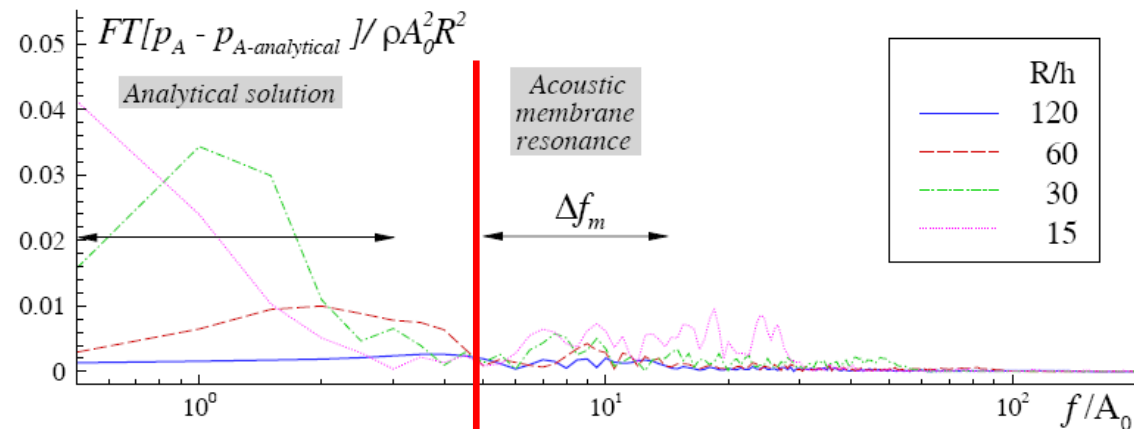
4. Stability? Pressure field quality?

Method 1 (artificial viscosity): Pressure field quality



Remember: we want pressure/force accuracy: we work for engineers, not gaming or movie people!

Where do the errors go? => acoustics



4. Stability? Pressure field quality?

Method 2

use of a **density diffusive term** (proportional to a Rusanov flux) in the continuity equation, e.g., δ -SPH, in addition to the artificial viscosity

$$\frac{d\rho_i}{dt} = -\rho_i \sum_{j \in \Omega} (\vec{v}_j - \vec{v}_i) \cdot \vec{\nabla}_i W_{ij} \omega_j + \delta h c_0 \sum_{j \in \Omega} \vec{\psi}_{ij} \cdot \vec{\nabla}_i W_{ij} \omega_j,$$

Method 3

use of **Riemann solvers** (Vila, 1999) between each pair of particles (standard in FVM for hyperbolic systems)

$$\left\{ \begin{array}{l} \frac{dx_i}{dt} \Big|_{v_0} = \bar{v}_{0i} \\ \frac{d\omega_i}{dt} \Big|_{v_0} = \sum_{j \in P(\Omega)} \omega_i \omega_j (\bar{v}_{0j} - \bar{v}_{0i}) \cdot \nabla_i W_{ij} \quad \text{given by the Riemann problem solution} \\ \frac{d\omega_i \bar{\phi}_i}{dt} \Big|_{v_0} + \sum_{j \in P(\Omega)} \omega_i \omega_j \left(\bar{F}_i - \bar{\phi}_i \otimes \bar{v}_{0i} + \bar{F}_j - \bar{\phi}_j \otimes \bar{v}_{0j} \right) \cdot \nabla_i W_{ij} = \omega_i \bar{S}_i \end{array} \right.$$

Method 4

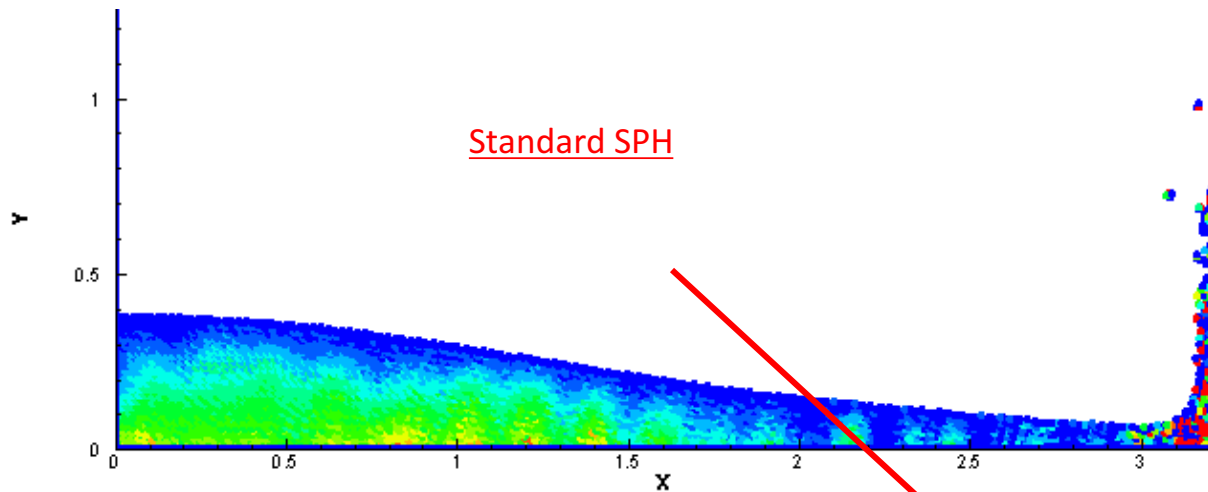
use an **incompressible** semi-implicit solution

Oger G. et al., SPH accuracy improvement through the combination of a quasi-Lagrangian shifting transport velocity and consistent ALE formalisms, J. Comput. Phys. **313**, 2016

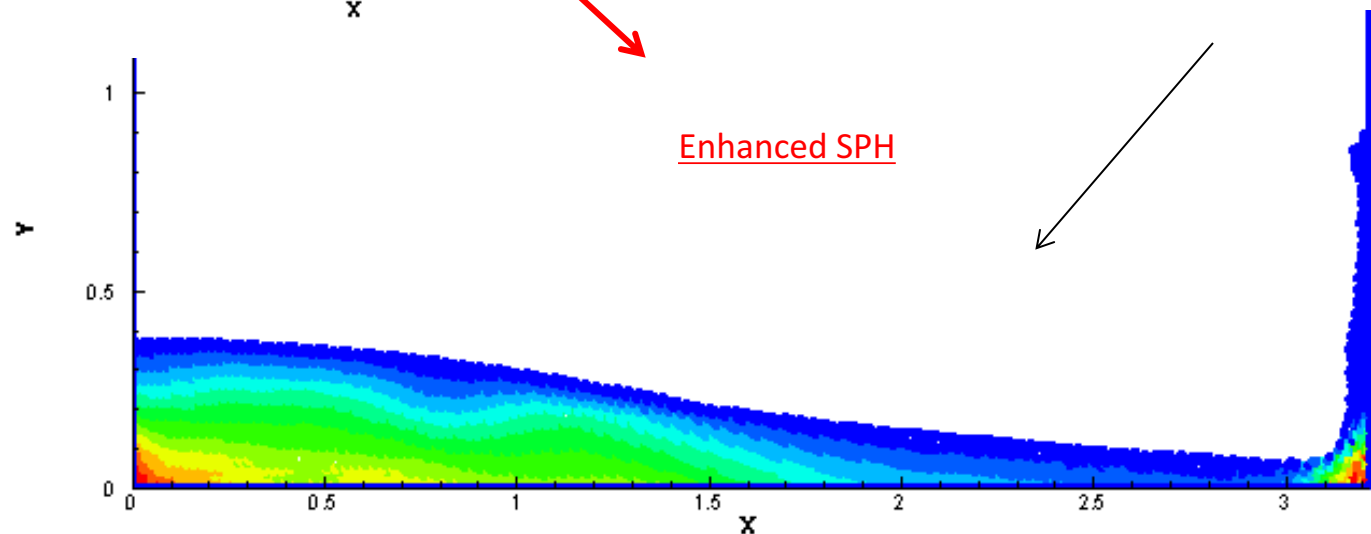
Antuono et al., Free-surface flows solved by means of SPH schemes with numerical diffusive terms, Comput. Phys. Commun. **181**, 2010

4. Stability? Pressure field quality?

Methods 2 to 4 : Pressure field quality



Use of Riemann solvers
or density diffusive term,
or incompressible variant



SPH fact checking



partly true

1. SPH cannot be solved implicitly (as mesh-based) / it has very small time steps

partly true

2. Weak-compressibility is an unphysical trick

not fully true

3. SPH conserves everything

no more true

4. SPH is unstable / SPH pressures are noisy

not so true

5. SPH is not convergent

not true

6. SPH is not accurate

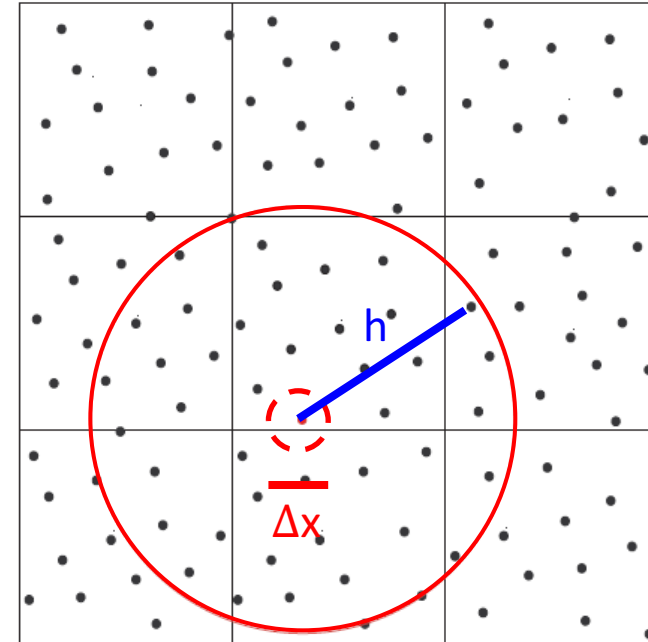
5. Convergence? 6. Accuracy?

A double convergence criteria

$$\langle u \rangle_i \cong \int u(\mathbf{x}) W(|\mathbf{x} - \mathbf{x}_i|) d\mathbf{x} \cong \sum_j u_j W(|\mathbf{x}_j - \mathbf{x}_i|) dV_j$$

$h \rightarrow 0$

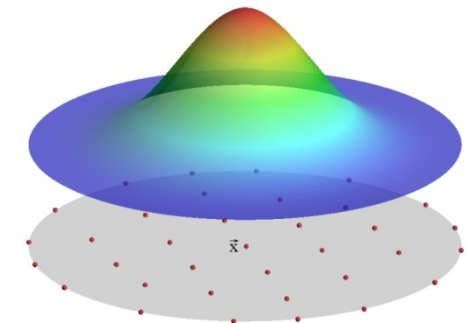
$\Delta x/h \rightarrow 0$



convergence order (Mas Gallic & Raviart):

$$h^2 + h^{-n} \left(\frac{\Delta x}{h} \right)^2$$

=> inconsistent if $\Delta x/h = \text{cst} !!!$
though common practice!

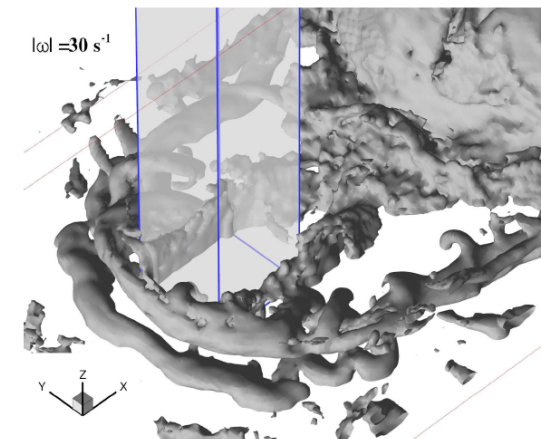
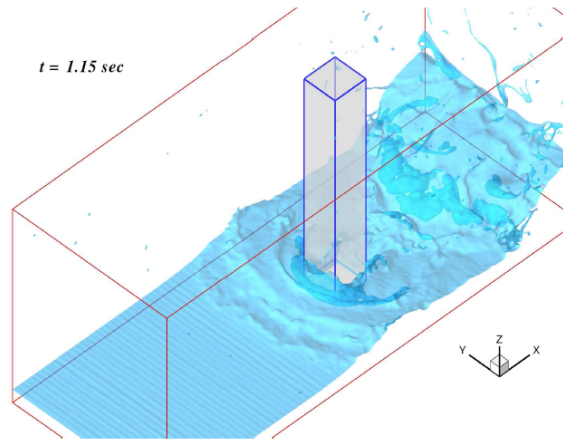
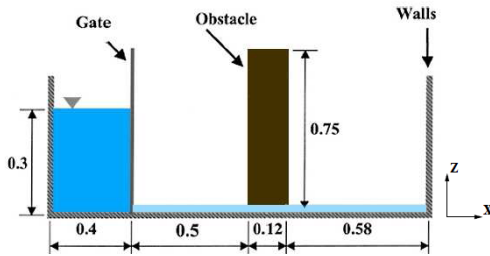


First-order operators (grad, div) used in standard SPH
diverge at order 1 (pressure gradient) or are not
convergent (velocity divergence)!!!

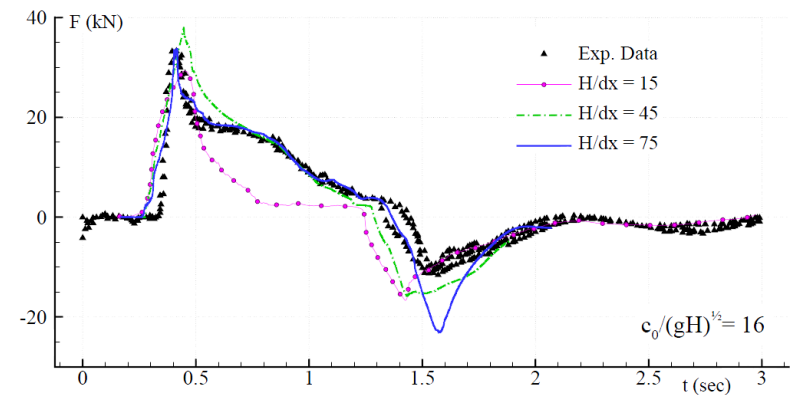
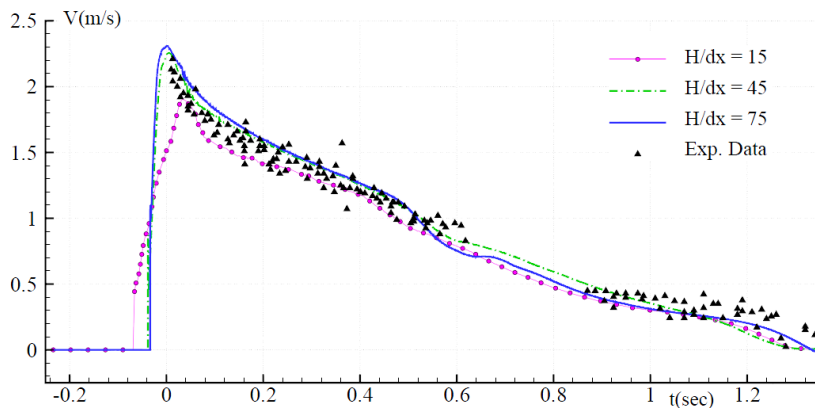
5. Convergence? 6. Accuracy?

And in practice, with $\Delta x/h = \text{cst}$?

Validation on dam breaking test cases



- Good prediction of velocity and forces:

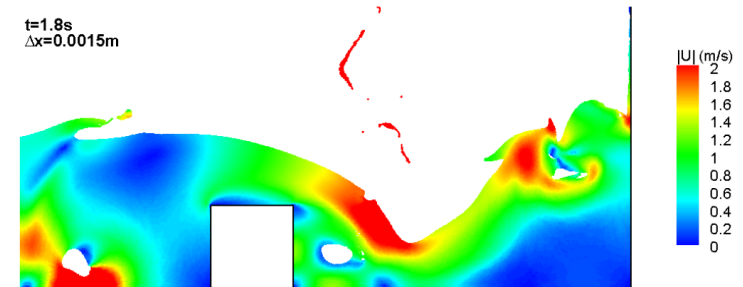
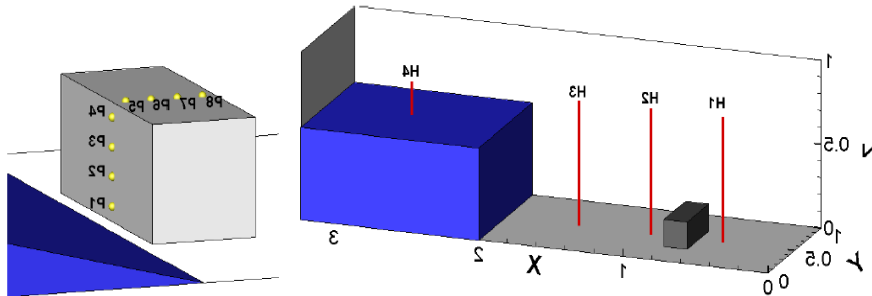


Marrone S. et al., δ -SPH model for simulating violent impact, *Comput. Meth. Appl. Mech. Engng.* **200**, 2011

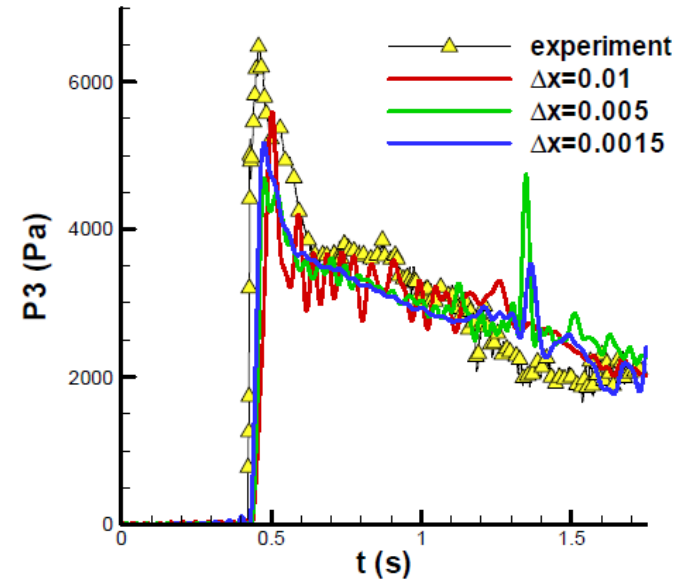
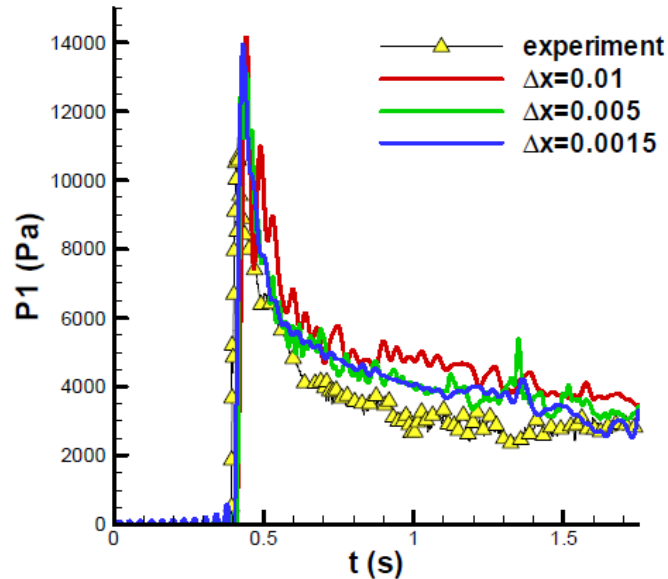
Marrone S. et al., Fast free-surface detection and level-set function definition in SPH solvers, *J. Comput. Phys.* **229**, 2010

5. Convergence? 6. Accuracy?

Validation on dam breaking test cases



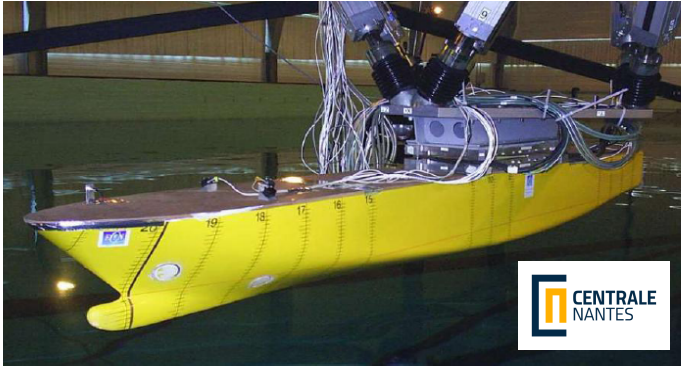
- Good prediction of pressures as well:



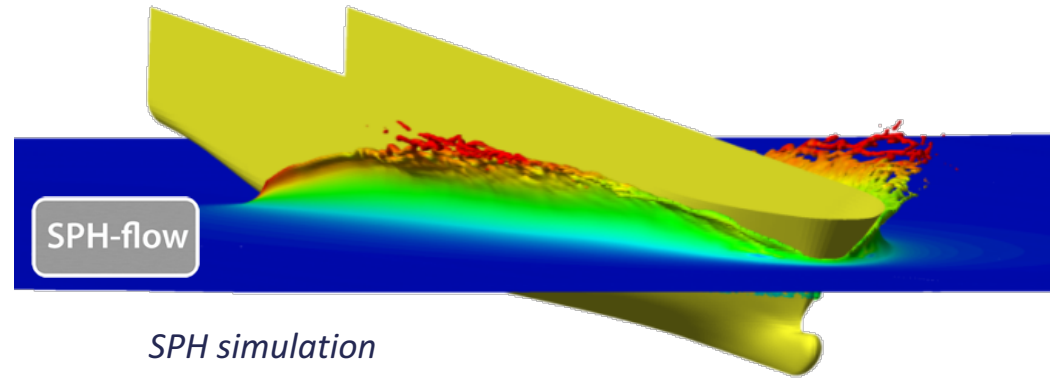
Heuristic convergence order is usually ~ 1

5. Convergence? 6. Accuracy?

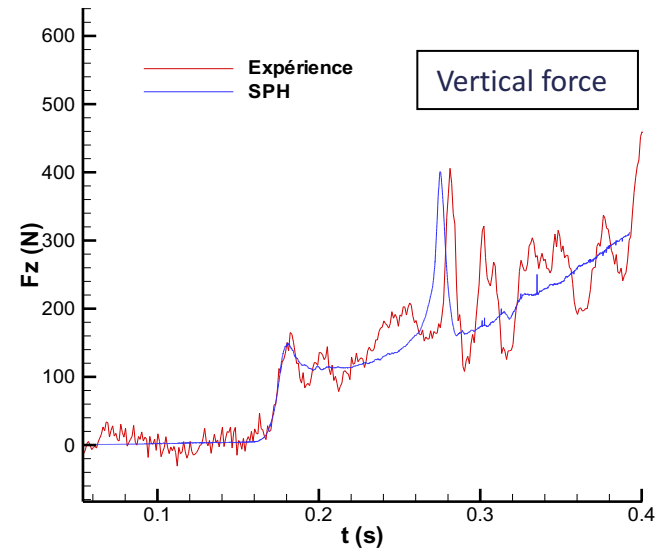
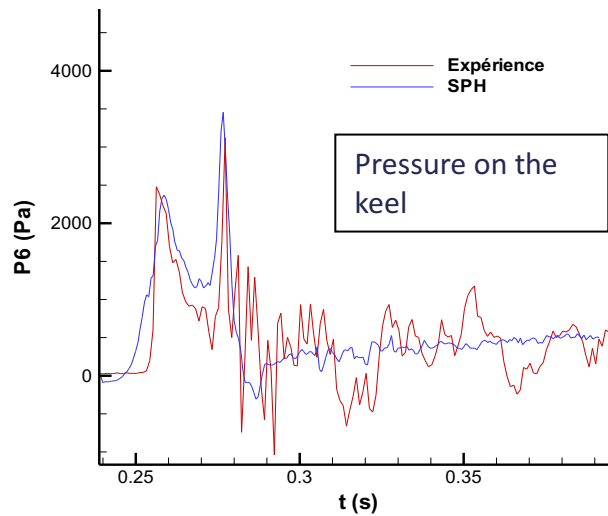
Validation on slamming impact on a real application



Experiment (ECN wave tank)



SPH simulation

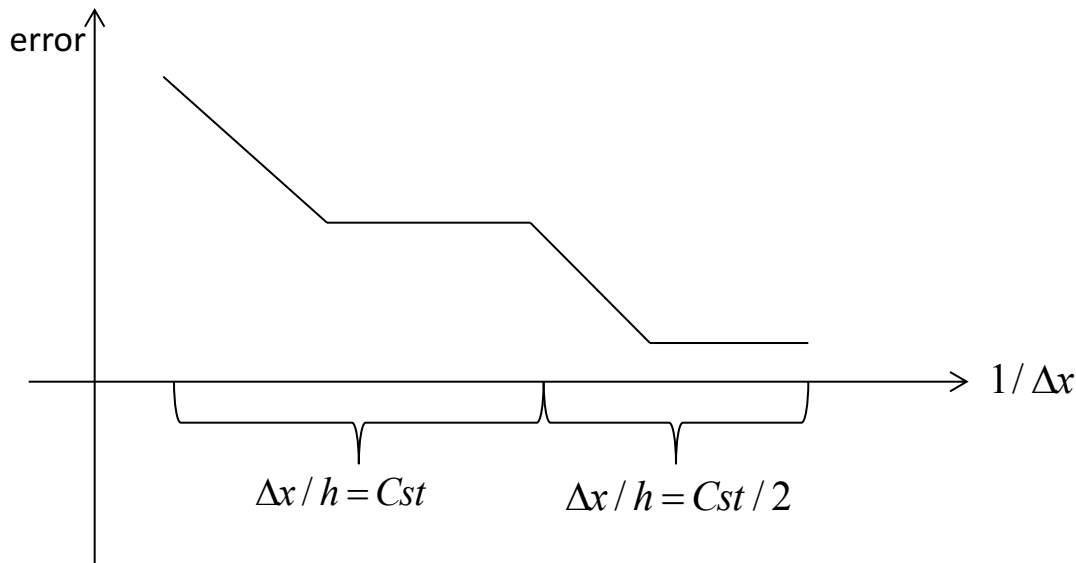


- 3 million particles
- 6m/s impact (real scale)
- 250m ship (real scale)

5. Convergence? 6. Accuracy?

So what can be the explanation?

- Lagrangian! = no discretisation of the convection term $\frac{d\mathbf{u}}{dt} = \frac{\partial \mathbf{u}}{\partial t} + (\text{grad} \mathbf{u})\mathbf{u}$
 => **exact convection**
 => provides accuracy for flows dominated by convection (**fast dynamics flows**), e.g. 3D dambreak w/ 80k particles
 => convergence/accuracy is a mixed between exact (convection) and poor (pressure gradient, velocity divergence)
- **Step-like convergence** (Quinlan, Ellero... *et al.*) + already **large stencil** (250 neighbors in 3D!)



⇒ Heuristic convergence often of order 1 with « tolerable » saturation (problem dependent!)

⇒ + « saved » by conservation and exact convection

- Discretization preserving **conservation** of mass, momenta and total energy

SPH fact checking



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5. SPH is not convergent

not true

6. SPH is not accurate

not yet!

7. SPH cannot be high order

false

8. SPH should be purely Lagrangian

7. Higher order? 8. Purely Lagrangian?

Operators can be easily corrected locally to increase their order of convergence (MLS...)

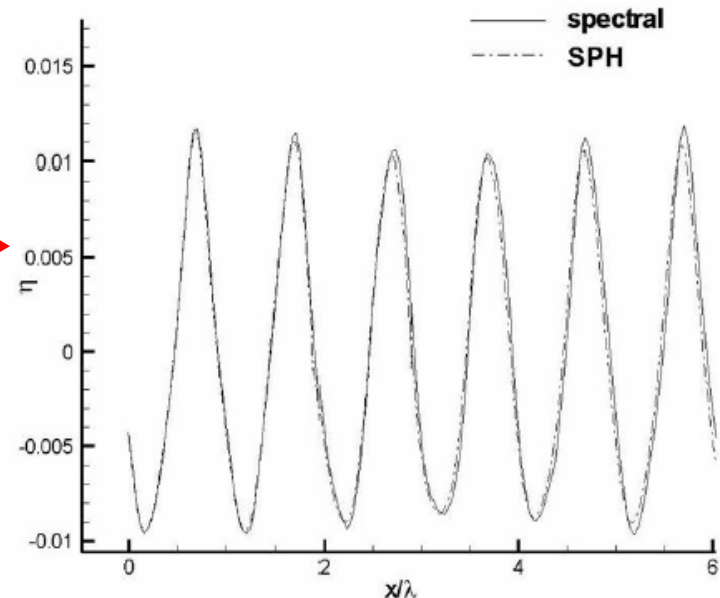
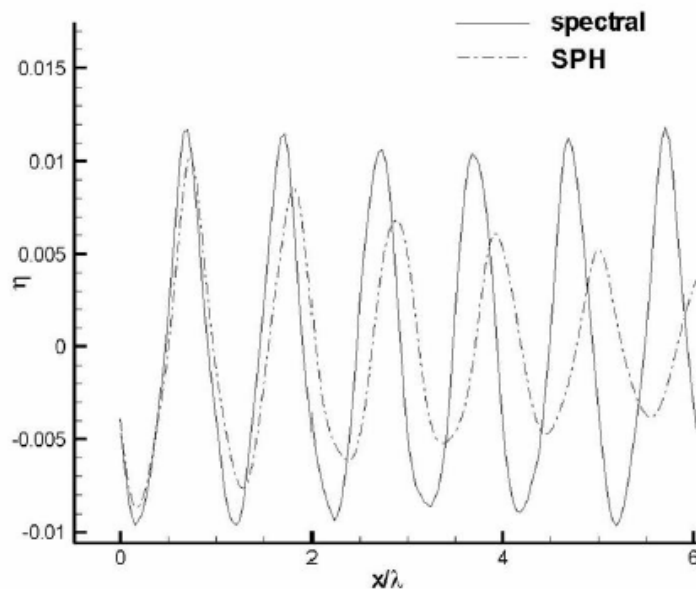
principle: imposing that the operator becomes exact for constant, linear, quadratic... fields

e.g.:

$$B_i = \left(\sum_{j \in \mathcal{P}} w_j (x_j - x_i) \otimes \nabla W_{ij} \right)^{-1}$$

$$\nabla W_{ij} \rightarrow \frac{1}{2} (B_i + B_j) \cdot \nabla W_{ij}$$

=> Efficiency on an academic example not meant to be solved by SPH: the propagation of a linear gravity wave (first-order convergent operators used)



7. Higher order? 8. Purely Lagrangian?

Consequences:

- => Convergence is theoretically **recstored** (see, e.g., Vila)
- => **Computational cost is increased** quite a lot (need to solve small matrices for each particle at each time step)
- => Typical corrections only restore order 1 (which was already heuristically obtained with reasonable saturation level) : order 2 is costly
 - => **more for accuracy than for convergence itself**
- => Correction impacts other aspects, **especially boundary conditions** => not so easy
- => **Calculations are often not improved/less stable!**

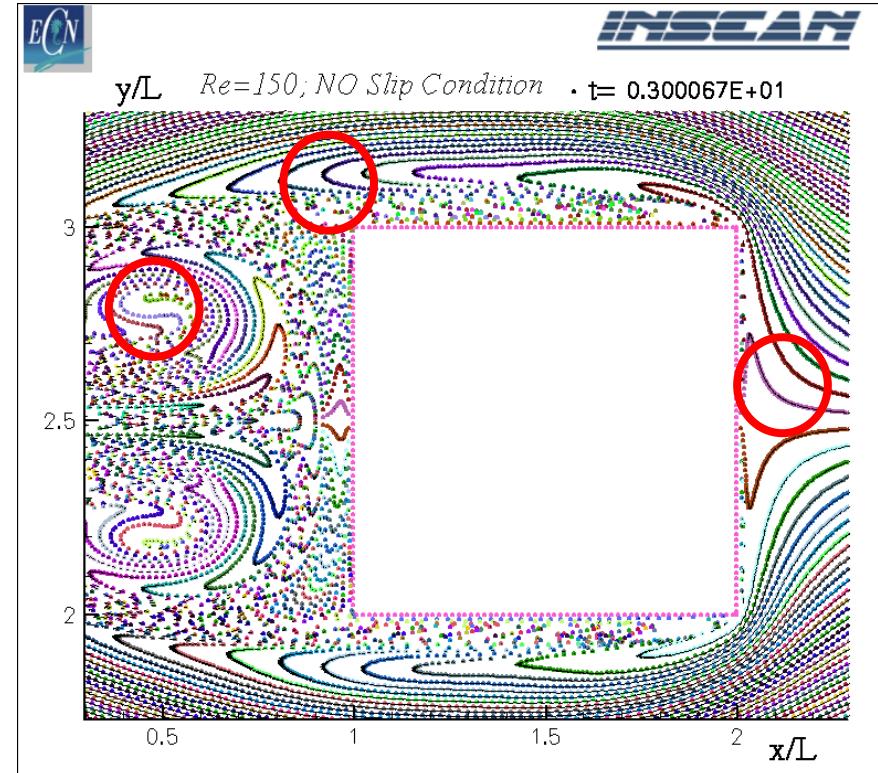
7. Higher order? 8. Purely Lagrangian?

More accuracy = TOO Lagrangian

- ⇒ harmful effects on spatial interpolation, accuracy and stability
- ⇒ leads to numerical artefacts (especially on pressure fields)
- ⇒ increases the numerical diffusion in the end

So, **the better the worse!** : the more accurate, the more Lagrangian and the more Lagrangian, the less accurate!

Together with difficulties at the boundaries, this explains why we do not see many high-order simulations...



- So why standard SPH particle distributions are so regular? => thanks to errors!
 - Hidden projection
 - Literature shows that errors on the pressure gradient induces a force tending to « fill voids »



7. Higher order? 8. Purely Lagrangian?

Solution found in recent years when the schemes became more accurate

⇒ People often use arbitrary shifting (XSPH, shifting...) to homogenize particle distribution in time
problem: it is not conservative and a lot comes from conservation!

⇒ We proposed a **fully-conservative** alternative through an **ALE formulation** (adapted from Vila, 1999 => Oger et al., J. Comput. Phys. **313**, 2016) used in a **quasi-Lagrangian** way (with same kind of displacements as shifting) => « **consistent shifting** »

$$\frac{d\vec{x}_i}{dt} = \vec{v}_{0i}, \quad (14)$$

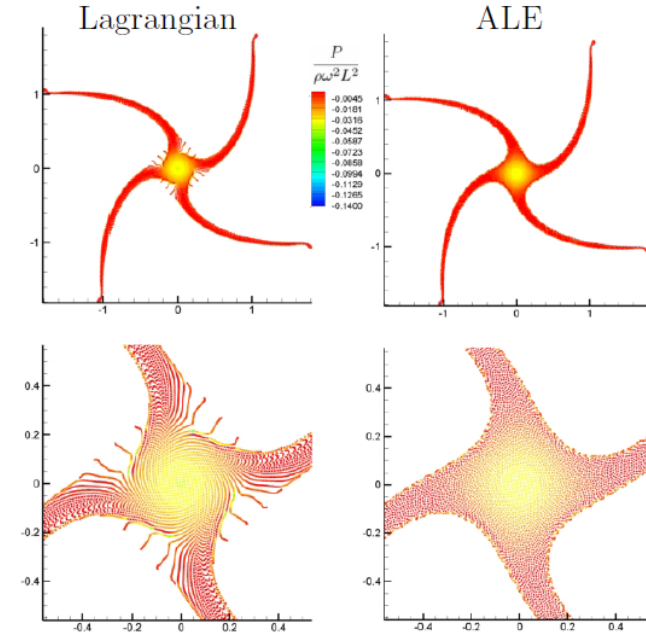
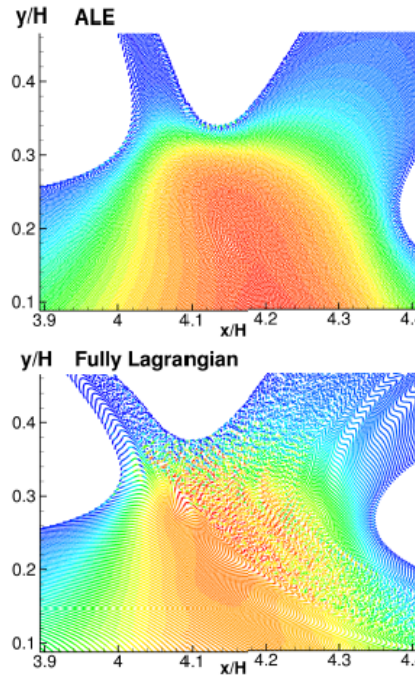
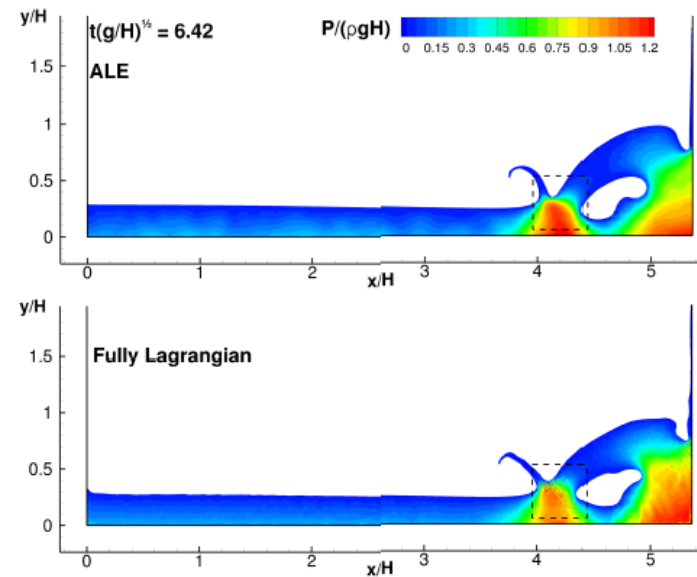
$$\frac{d\omega_i}{dt} = \omega_i \sum_j (\vec{v}_{0j} - \vec{v}_{0i}) \nabla W_{ij} \omega_j, \quad (15)$$

$$\frac{d(\omega_i \rho_i)}{dt} = -\omega_i \sum_j (\rho_i (\vec{v}_i - \vec{v}_{0i}) + \rho_j (\vec{v}_j - \vec{v}_{0j})) \nabla W_{ij} \omega_j, \quad (16)$$

$$\begin{aligned} \frac{d(\omega_i \rho_i \vec{v}_i)}{dt} = & -\omega_i \sum_j \left(\rho_i \vec{v}_i \otimes (\vec{v}_i - \vec{v}_{0i}) + \rho_j \vec{v}_j \otimes (\vec{v}_j - \vec{v}_{0j}) + P_i \bar{\bar{I}} + P_j \bar{\bar{I}} \right) \nabla W_{ij} \omega_j \\ & + \omega_i \rho_i \vec{g}. \end{aligned} \quad (17)$$

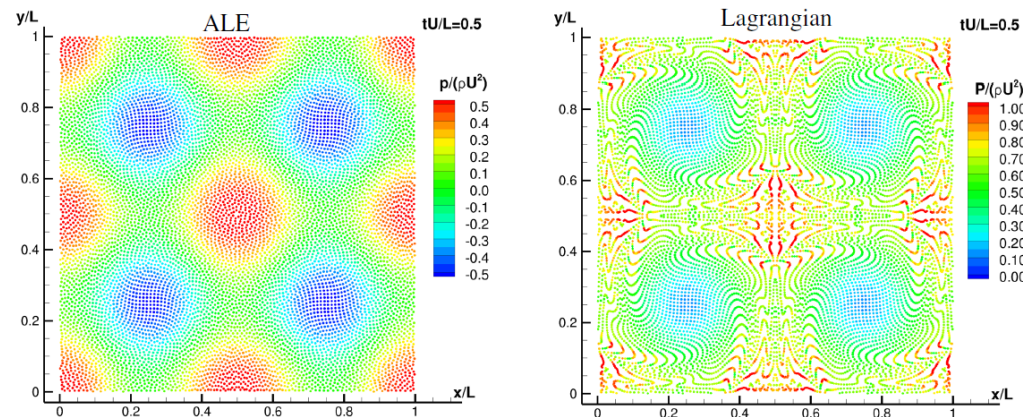
7. Higher order? 8. Purely Lagrangian?

Examples of results with ALE formulation and Riemann solvers



With this ALE formulation we have all the ingredients to build a higher-order model (work in progress)

Oger G. et al., SPH accuracy improvement through the combination of a quasi-Lagrangian shifting transport velocity and consistent ALE formalisms, *J. Comput. Phys.* **313**, 2016



SPH fact checking



partly true

1. SPH cannot be solved implicitly (as mesh-based) / it has very small time steps

partly true

2. Weak-compressibility is an unphysical trick

not fully true

3. SPH conserves everything

no more true

4. SPH is unstable / SPH pressures are noisy

not so true

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false

8. SPH should be purely Lagrangian

mainly false

9. Free-surface conditions are not modelled whereas they should

9.-10. Boundary conditions

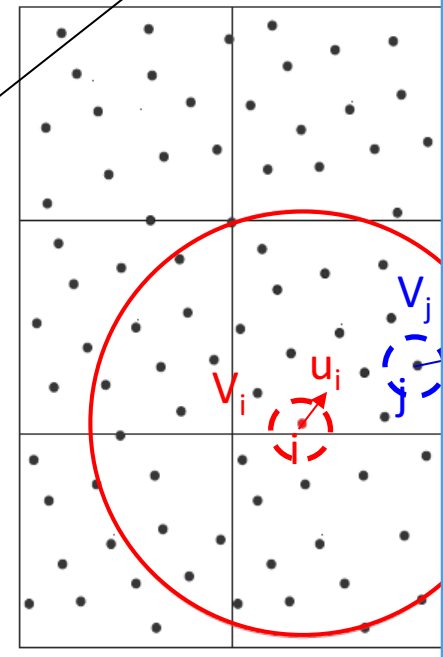
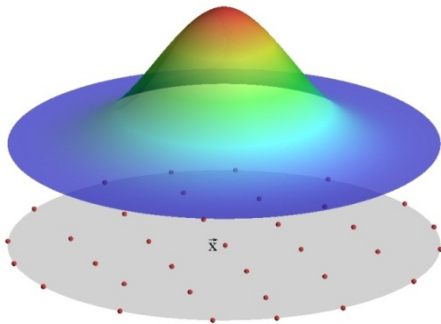
=> mollification + discretisation

$$\langle \mathbf{u} \rangle_i = \int \mathbf{u}(\mathbf{x}) \delta(|\mathbf{x} - \mathbf{x}_i|) d\mathbf{x} \stackrel{\textcircled{1}}{\cong} \int \mathbf{u}(\mathbf{x}) W(|\mathbf{x} - \mathbf{x}_i|) d\mathbf{x} \stackrel{\textcircled{2}}{\cong} \sum_j \mathbf{u}_j W(|\mathbf{x}_j - \mathbf{x}_i|) V_j$$

Known at the neighbour scattered points (particle)

Volume associated to the scattered points (particles)

Analytical function (kernel)



Incomplete support

No more true

+ transfer of the differentiation to the kernel to get differential operators:

$$\nabla f(\vec{r}) \stackrel{\textcircled{1}}{\approx} \int_D \nabla f(\vec{x}) W(\vec{r} - \vec{x}) d\vec{x} = \int_D f(\vec{x}) \nabla W(\vec{r} - \vec{x}) d\vec{x}$$

analytical

9.-10. Boundary conditions



=> **boundary terms** need to be accounted for in integration by parts

+

the support is no more filled by neighbours close to the boundary => potential inaccuracies

9. Free-surface boundary conditions

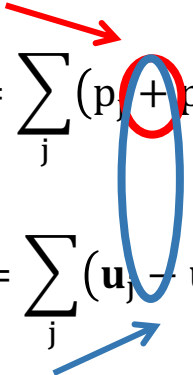
Kinematic FSBC is straightforward since we have a Lagrangian formalism


=> ok at convergence even though no particle strictly lies on the free surface due to their volume

We have proved that **dynamic FSBC** is verified in an integral sense provided appropriate operators are used for the pressure gradient, velocity divergence and velocity Laplacian (Colagrossi et al., Phys. Rev. E 2009 et 2011)

=> same reasoning as before but with boundary terms

for **Momentum** conservation (action/reaction)

$$\langle \nabla p \rangle_i = \sum_j (p_j + p_i) W(|x_j - x_i|) V_j$$


$$\langle \text{div} \mathbf{u} \rangle_i = \sum_j (\mathbf{u}_j - \mathbf{u}_i) \cdot W(|x_j - x_i|) V_j$$


for **Energy** conservation (Hamiltonian)

Colagrossi A. et al., Theoretical considerations on the free-surface role in the smoothed-particle-hydrodynamics model, Phys. Rev. E **79**, 2009

Colagrossi A. et al., Theoretical analysis and numerical verification of the consistency of viscous smoothed-particle-hydrodynamics formulations in simulating free-surface flows, Phys. Rev. E **84**, 2011

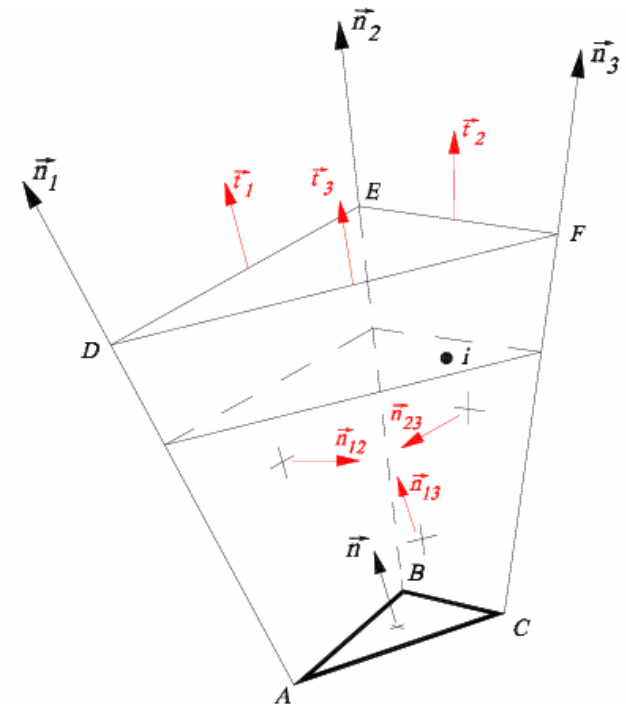
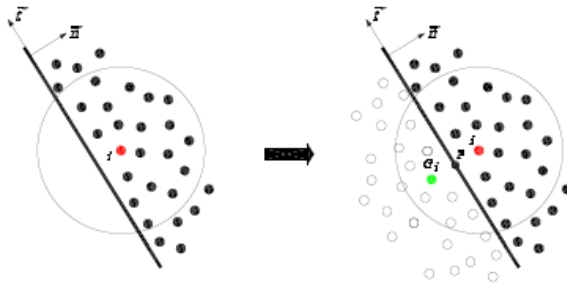
SPH fact checking



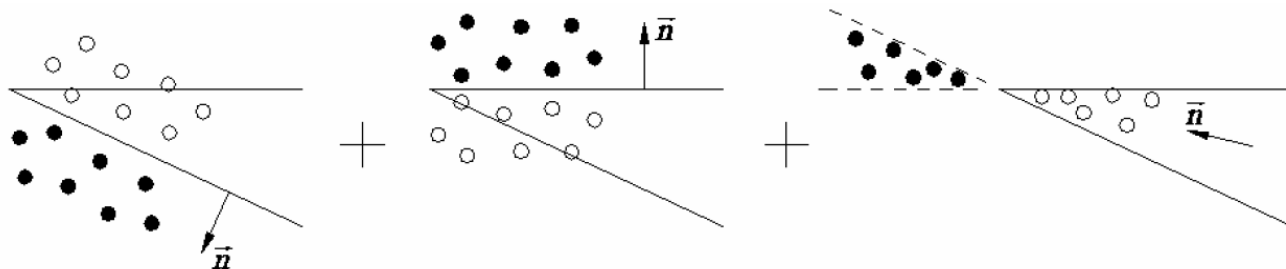
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| not so true | 5. SPH is not convergent |
| not true | 6. SPH is not accurate |
| not yet! | 7. SPH cannot be high order |
| false | 8. SPH should be purely Lagrangian |
| mainly false | 9. Free-surface conditions are not modelled whereas they should |
| no more true | 10. There is no good scheme to model wall boundary conditions |

10. Good wall boundary conditions?

Formerly we used the ghost method, best compromise between generality and accuracy
 => 3D generic technique from any surface (e.g. IGES format)

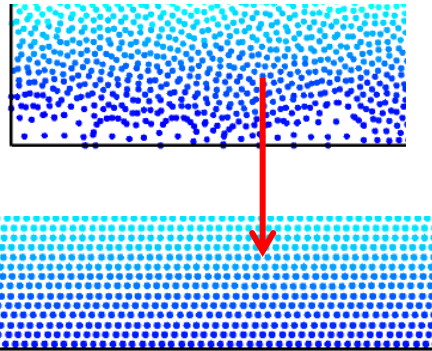


- Exact only for a flat panel, **difficulties with geometrical singularities**, especially sharp edges



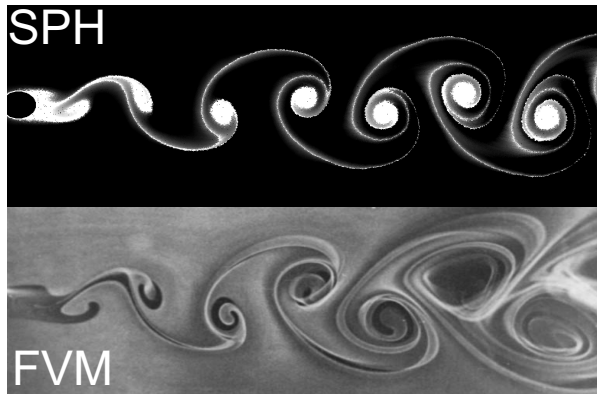
10. Good wall boundary conditions?

We proved that the no-slip ghost condition should not be applied to the velocity divergence to preserve the hyperbolicity of the system inviscid part (De Lefle et al., 6th SPHERIC workshop, 2011)



$$\begin{aligned} \frac{dE}{dt} &= - \sum_{i \in P(\Omega)} \sum_{j \in P(\Omega) \cup P(\partial\Omega)} m_i m_j \left(\frac{p_j \bar{v}_i}{\rho_j^2} + \frac{p_i \bar{v}_j}{\rho_i^2} + \frac{1}{2} \Pi_{ij} (\bar{v}_i + \bar{v}_j) \right) \cdot \nabla_i W_{ij} \\ &= \Delta E + \Delta E^\Pi \end{aligned}$$

Validation on the flow past cylinder at Re=200



	Strouhal	C _d
SPH	2.0	1.47
Experiment	1.9	1.3

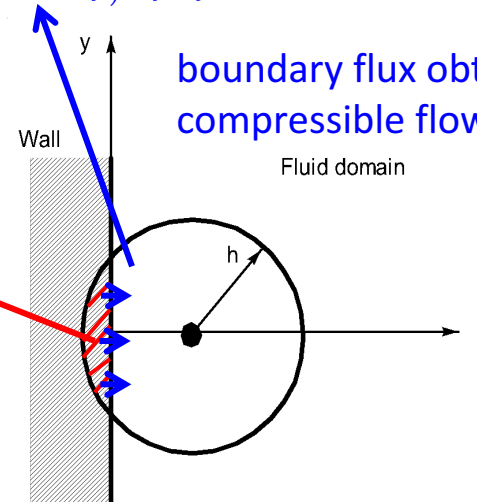
10. Good wall boundary conditions?

We proposed the **Normal Flux Method** (submitted to JCP)

$$\begin{cases} \left. \frac{dx_i}{dt} \right|_{v_0} = \overline{v_{0i}} \\ \left. \frac{d\omega_i}{dt} \right|_{v^0} = \frac{1}{\gamma_i} \sum_{j \in P(\Omega)} \omega_i \omega_j (\overline{v_{0j}} - \overline{v_{0i}}) \cdot \nabla_i W_{ij} + \frac{1}{\gamma_i} \sum_{j \in P(\partial\Omega)} \omega_i s_j (\overline{v_{0j}} - \overline{v_{0i}}) \cdot \overline{n_j} W_{ij} \\ \left. \frac{d\omega_i \overline{\phi_i}}{dt} \right|_{v^0} + \frac{1}{\gamma_i} \sum_{j \in P(\Omega)} \omega_i \omega_j (\overline{F_i} + \overline{F_j}) \cdot \nabla_i W_{ij} + \frac{1}{\gamma_i} \sum_{j \in P(\partial\Omega)} \omega_i s_j (\overline{F_i} + \overline{F_j}) \cdot \overline{n_j} W_{ij} = \omega_i \overline{S_i} \end{cases}$$

compensates missing
part of the kernel
support

boundary flux obtained from
compressible flow characteristics



- Belongs to the family of boundary integration techniques (like, e.g., USAW)
- Fully general technique
- No leakage of particles
- Permits very complex geometrical configurations meshed with millions of elements

SPH fact checking

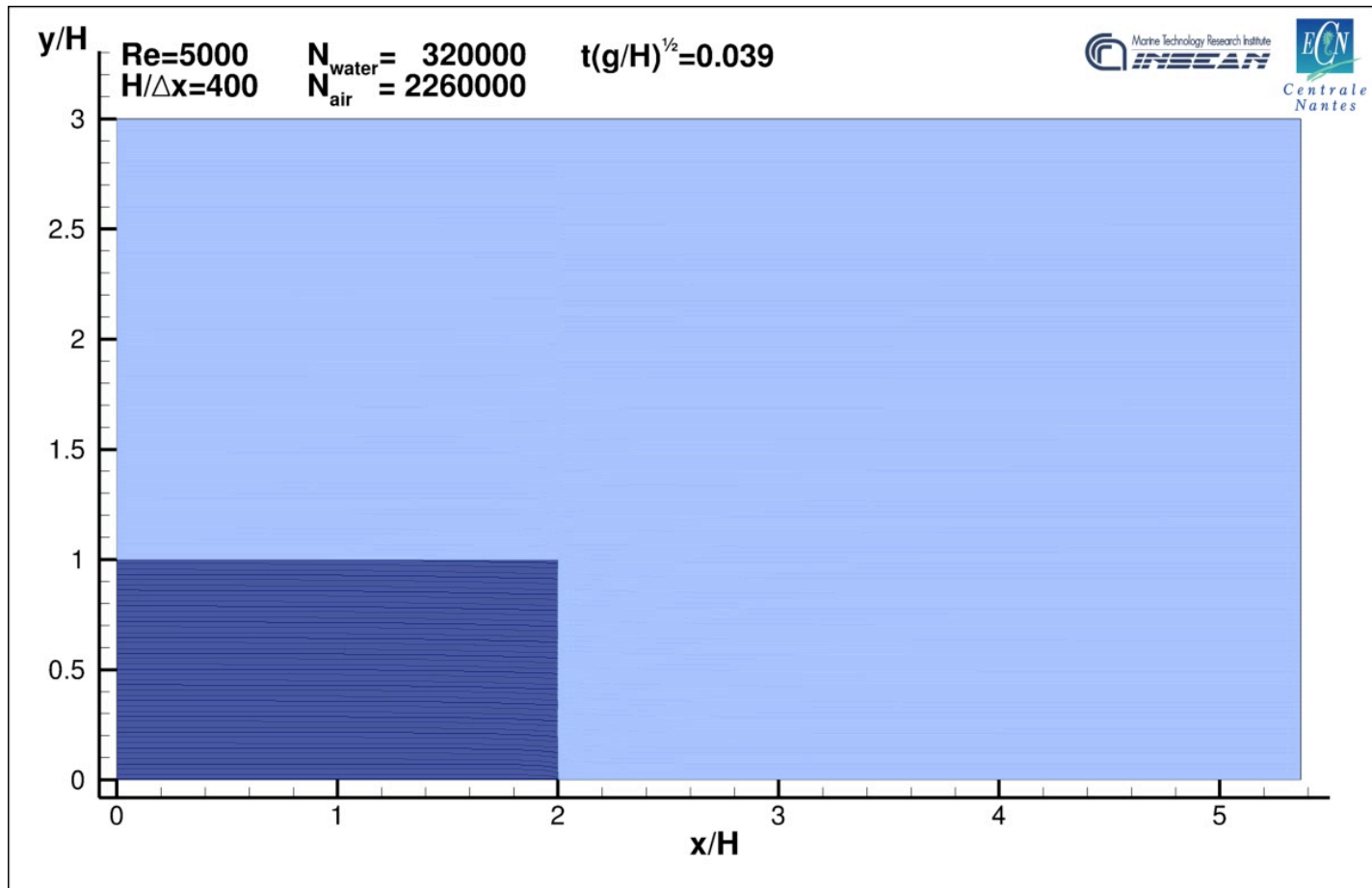


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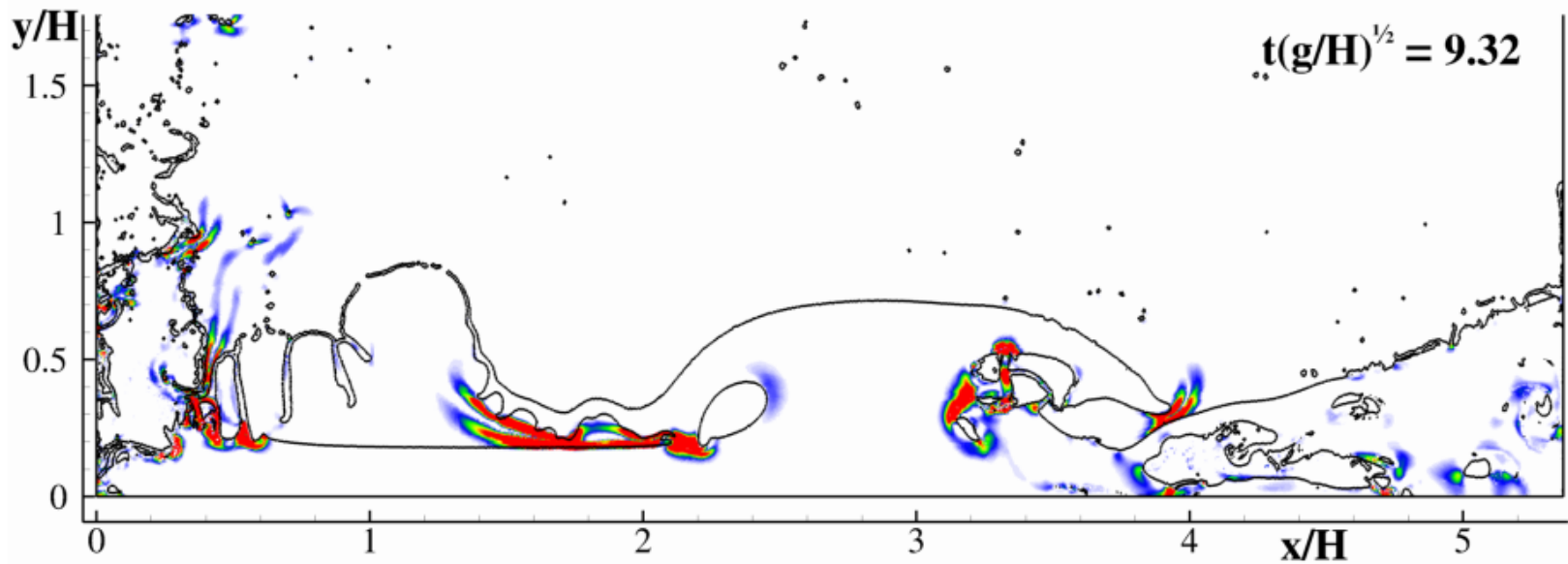
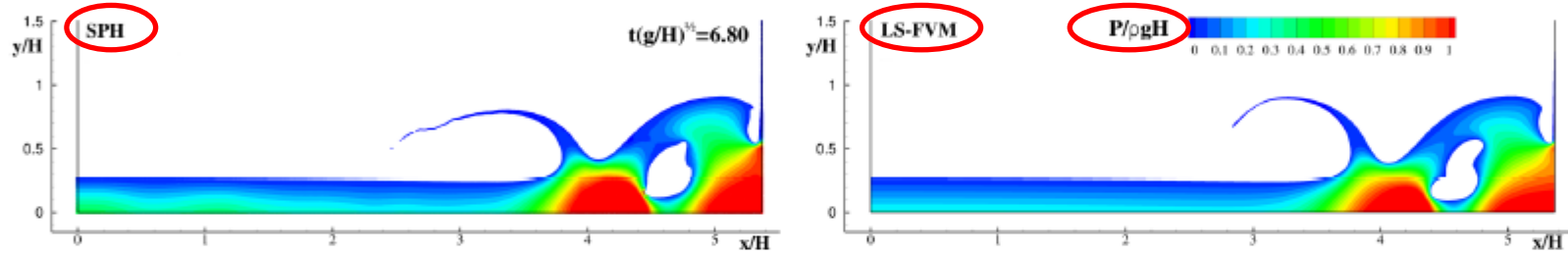
11. Validity of single phase approximation?

Lots of SPH users do not pay attention to the fact that **single-phase simulation is a priori limited to non-breaking free-surface flows**

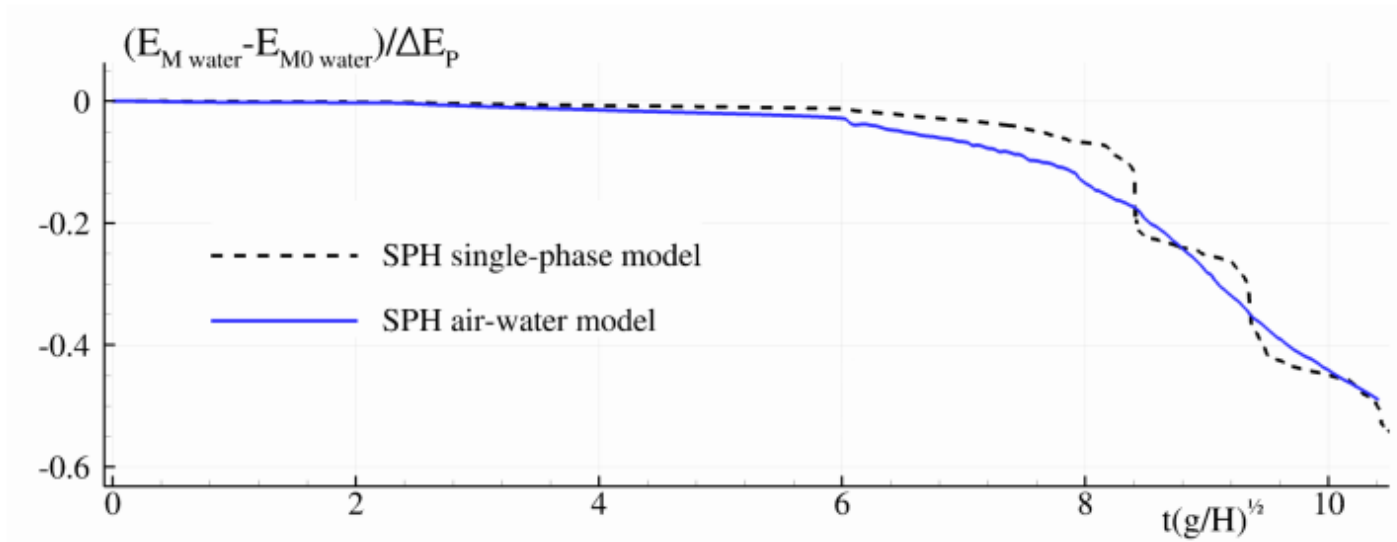
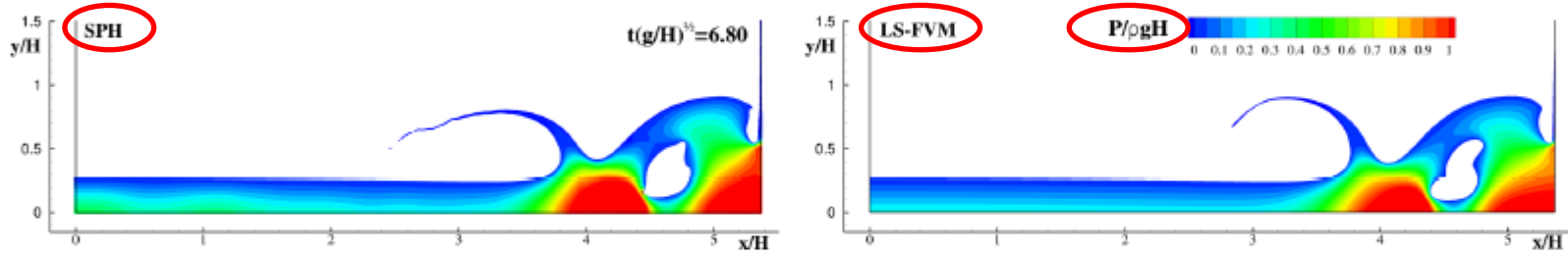
=> Study of a complex dambreak flow with 1 and 2 phases simulated



11. Validity of single phase approximation?



11. Validity of single phase approximation?



SPH fact checking

partly true

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mainly false

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11. Single-phase assumption used in free-surface SPH is physically a non-sense

still true

12. SPH is costly / any future tracks?

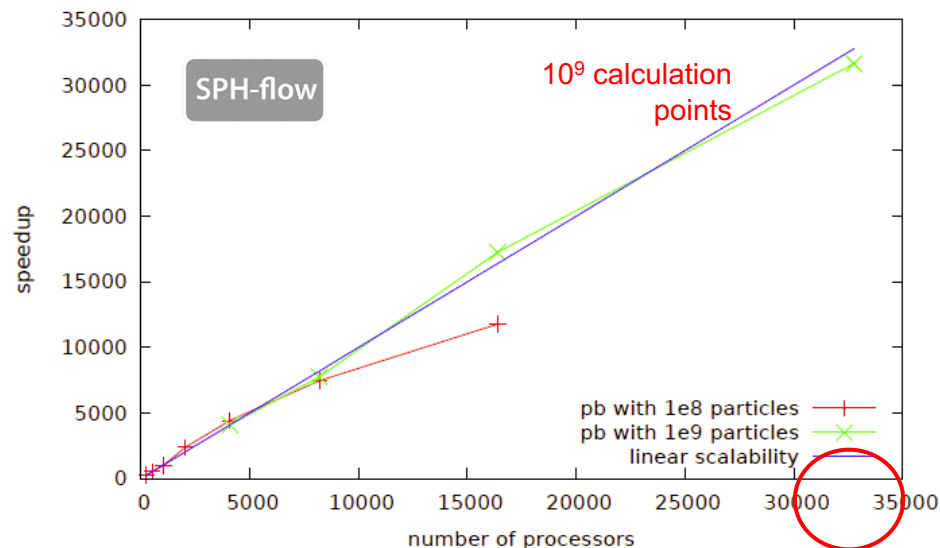
12. SPH is costly?

Unfortunately : yes!

- **Very large stencils** (typically 50 neighbours in 2D and 250 in 3D) => maybe less when we will have a robust second-order scheme?
- **Small time steps** => maybe a fully-implicit variant some day?

⇒ Compete well only **where mesh-based methods have difficulties** and for **fast dynamic flows**

⇒ Need for large hardware/efficient strategy, e.g. **MPI/OpenMP** or **GPGPU**



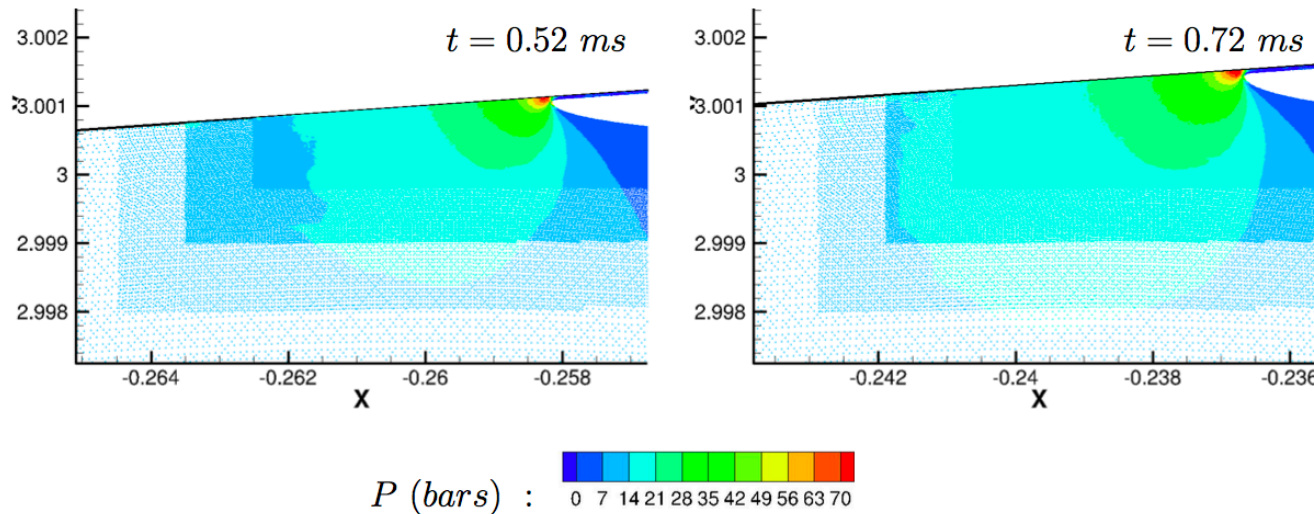
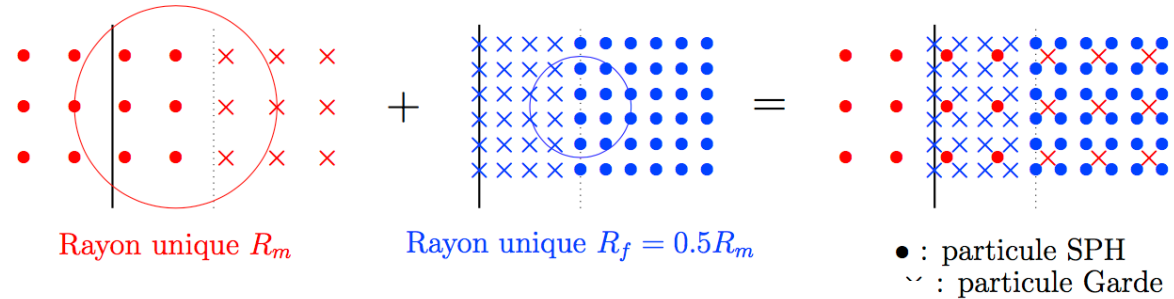
Linear scalability of SPH-Flow

32000
cores

12. Future tracks

Adaptive particle refinement (APR)

- Inspired from mesh-based AMR, but adapted to a Lagrangian formalism => use of « guard particles », prolongations, restrictions. as in AMR
- Has proved to be accurate, efficient and robust

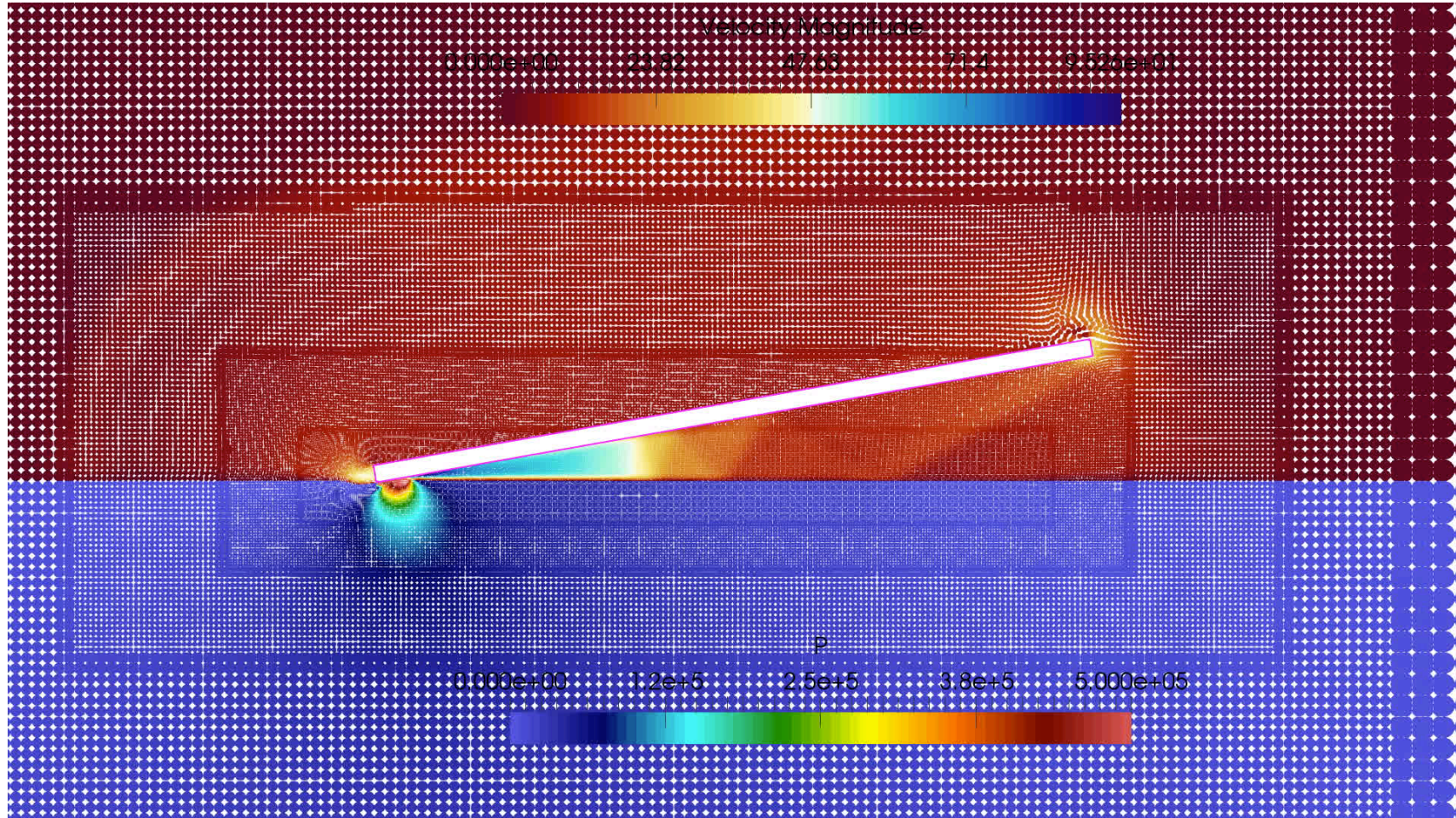


Barcarolo et al., Adaptive particle refinement and derefinement applied to the smoothed particle hydrodynamics method, *J. Comput. Phys.* **273**, 2014

Chiron et al., Analysis and improvements of Adaptive Particle Refinement (APR) through CPU time, accuracy and robustness considerations, to appear in *J. Comput. Phys.*, 2017

12. Future tracks

Adaptive particle refinement (APR): 40 m/s plate ditching (two-phase model)



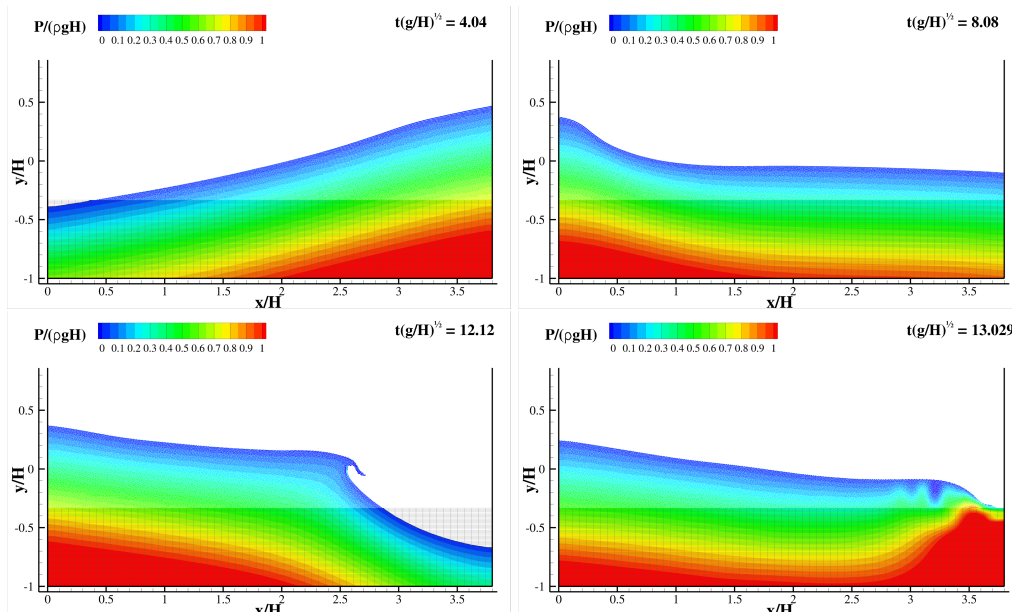
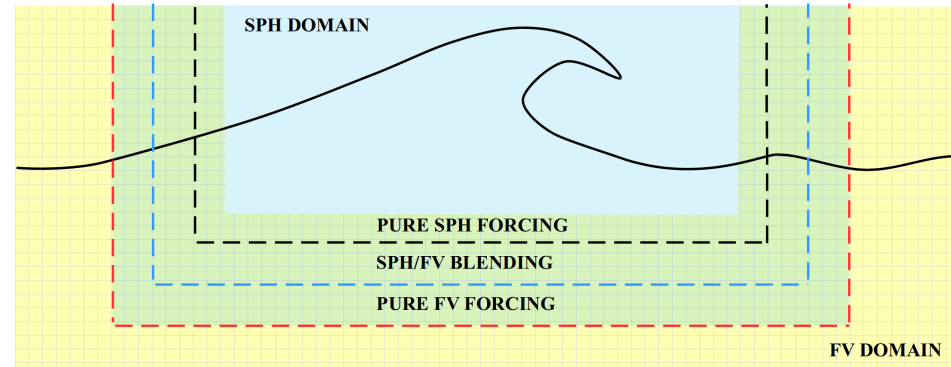
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12. Future tracks

Coupling to Finite Volumes

- Coupling with a **finite volume level-set solver**
- Principle : use of **forcing and blending zones**
- Efficient even when the change of solver intersects the free surface
- Validated on numerous 2D test cases



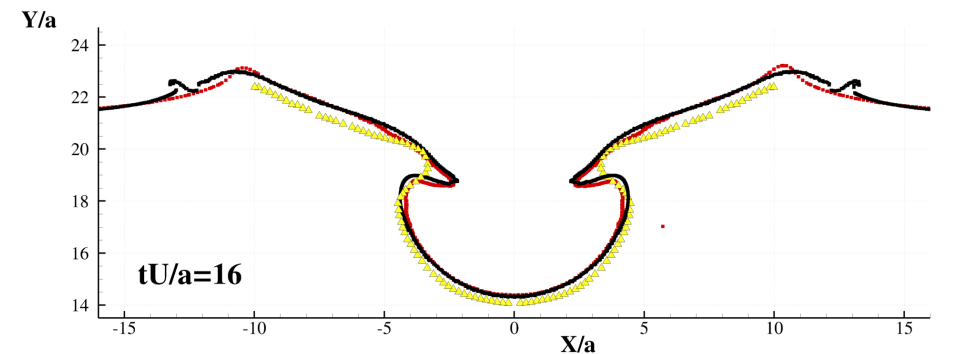
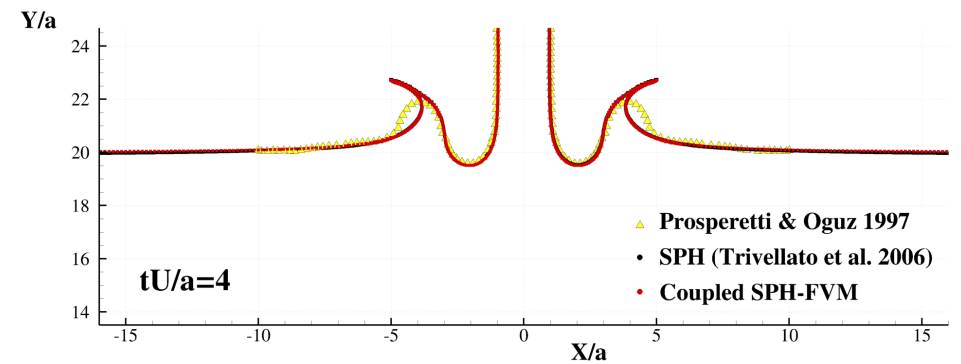
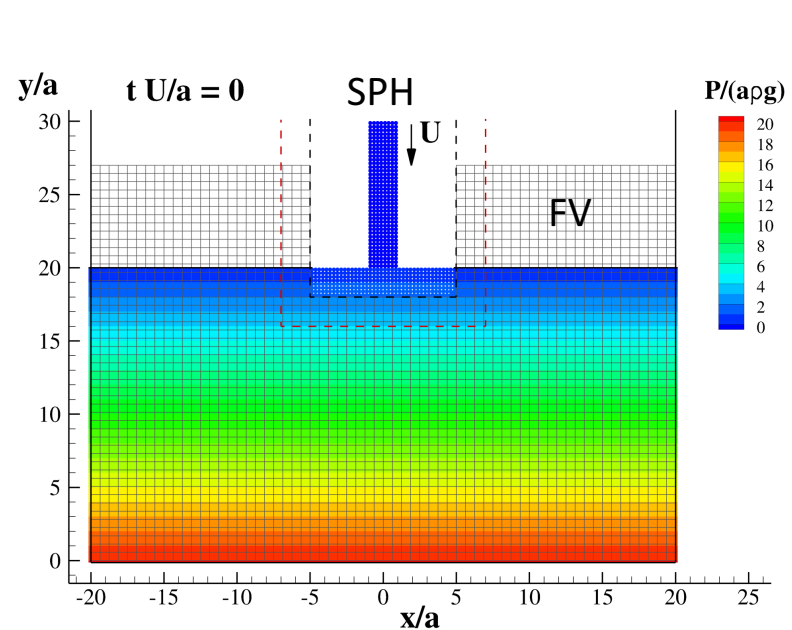
Marrone S. et al., Coupling of Smoothed Particle Hydrodynamics with Finite Volume method for free-surface flows, J. Comput. Phys. **310**, 2016

Marrone S. et al., Coupled SPH-FV method with net vorticity and mass transfer, submitted to J. Comput. Phys.

12. Future tracks

Coupling to Finite Volumes

Free-surface (Froude) driven test-case (difficult for the FV level-set solver)



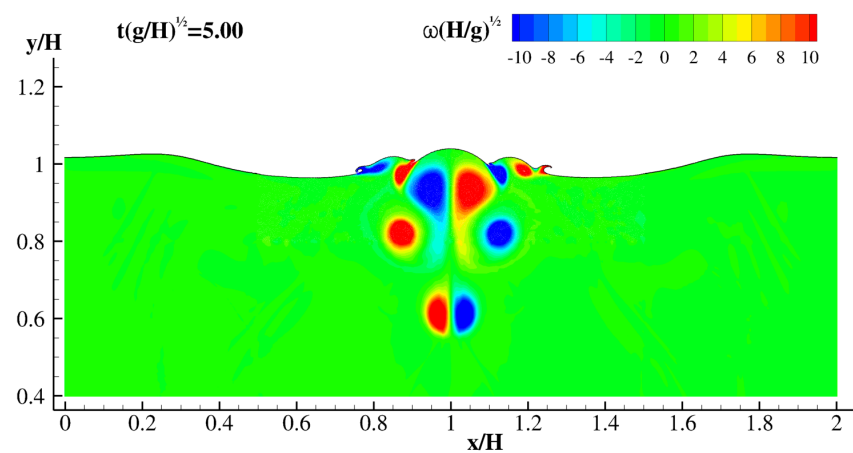
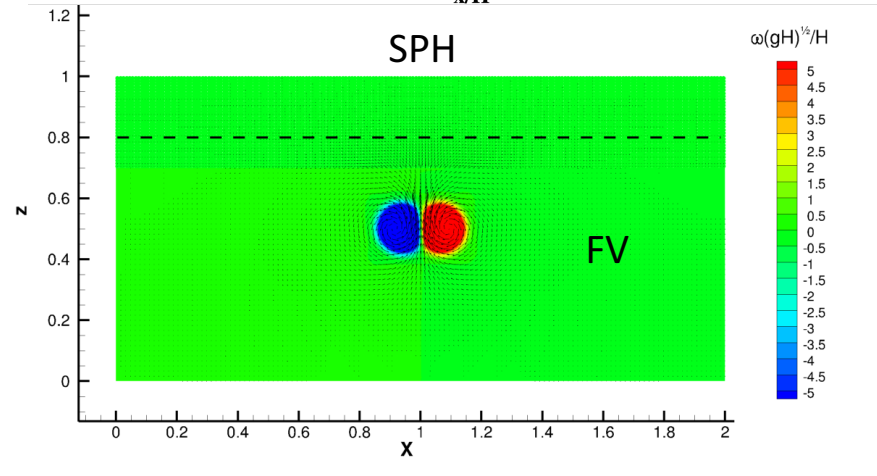
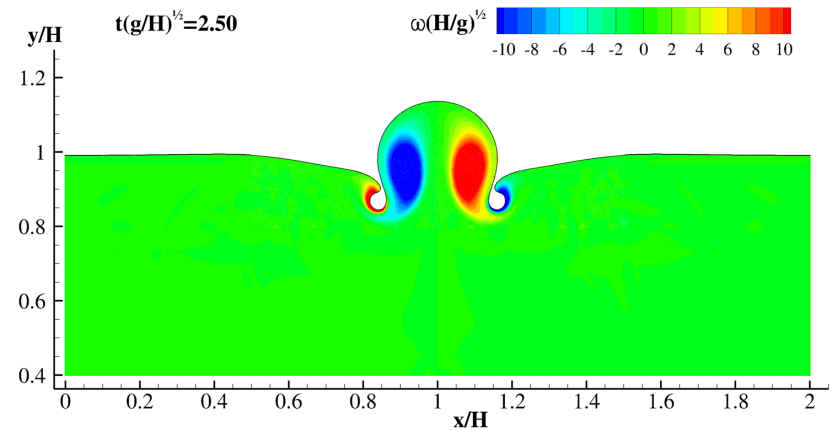
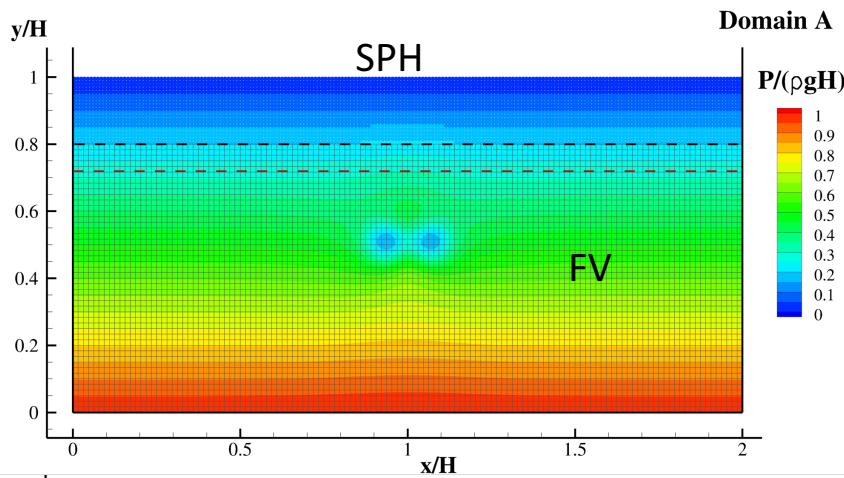
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12. Future tracks

Coupling to Finite Volumes

Vorticity-driven (Reynolds) test-case (difficult for the SPH solver)



Marrone S. et al., Coupling of Smoothed Particle Hydrodynamics with Finite Volume method for free-surface flows, J. Comput. Phys. **310**, 2016

Marrone S. et al., Coupled SPH-FV method with net vorticity and mass transfer, submitted to J. Comput. Phys.

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| still true | 12. SPH is costly / any future tracks? |
| no more true | 13. SPH is a research object which has no industrial potential |

13. Industrial applications?



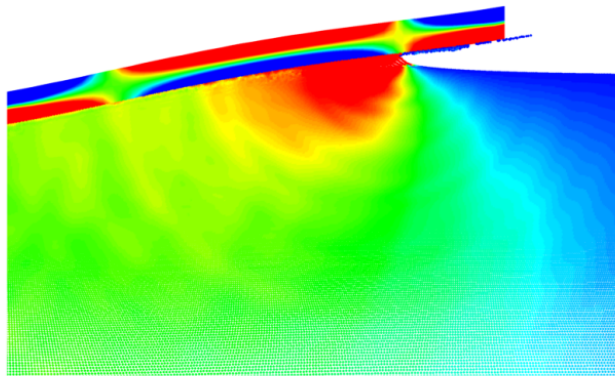
What is true

- SPH is **costly** and **not accurate for all problems**
=> it has a restricted field of applications (for now and probably for a long time)
- The fields of applications are:
 - **Fast dynamics problems**: small time steps + exact convection + complex interfaces
 - **Complex physics**: multi-body in the flow with contact, problems with different species/physical phenomena (explicit solving)
 - **Multi-solver problems**: easy coupling with other solvers: SPH-FEM / SPH-DEM / ...
- Another asset: **CAO to CFD** (like LBM)

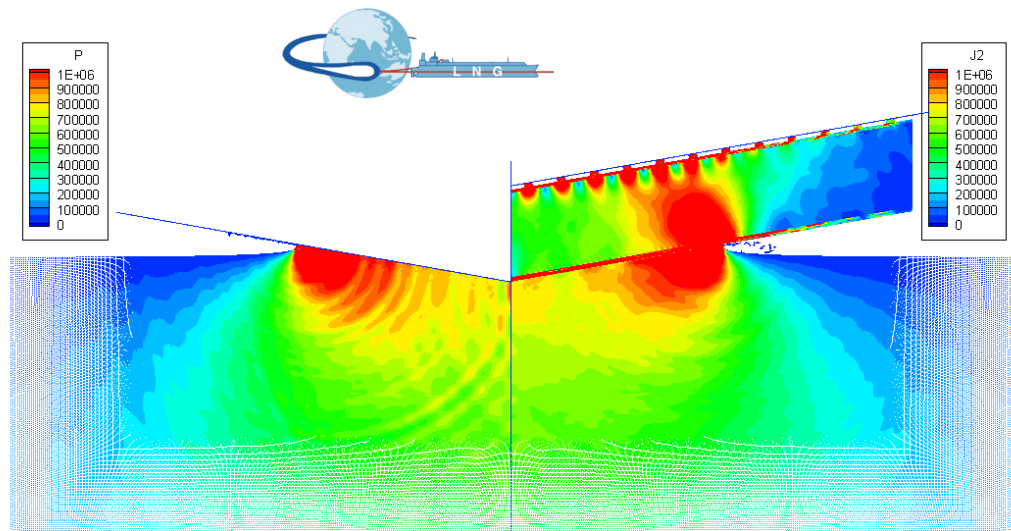
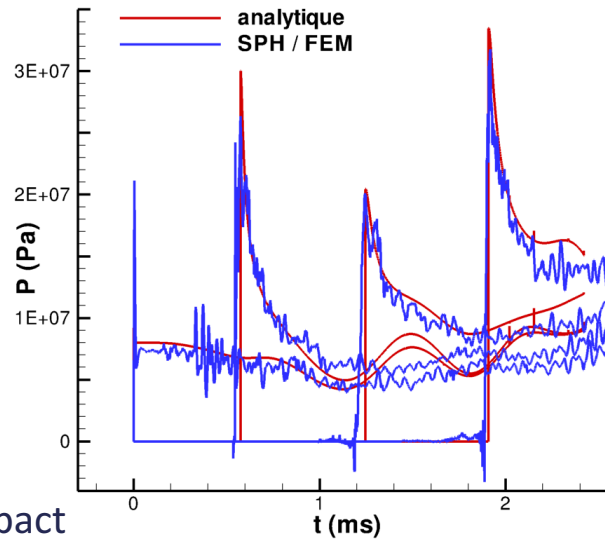
A rather extensive review: Shadloo et al., Computers & Fluids **136**, 2016

13. Industrial applications?

Fast fluid-structure coupled impact



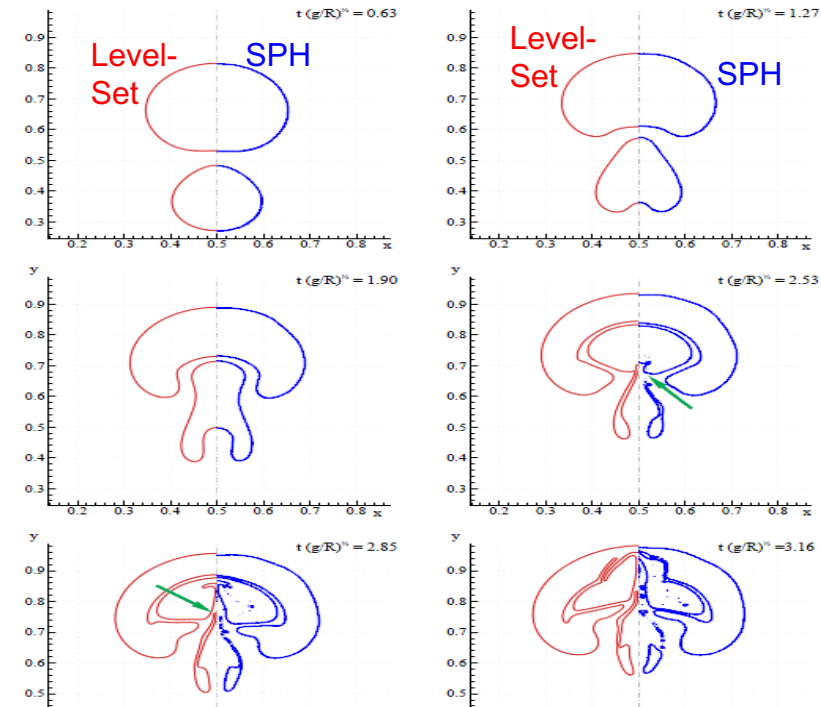
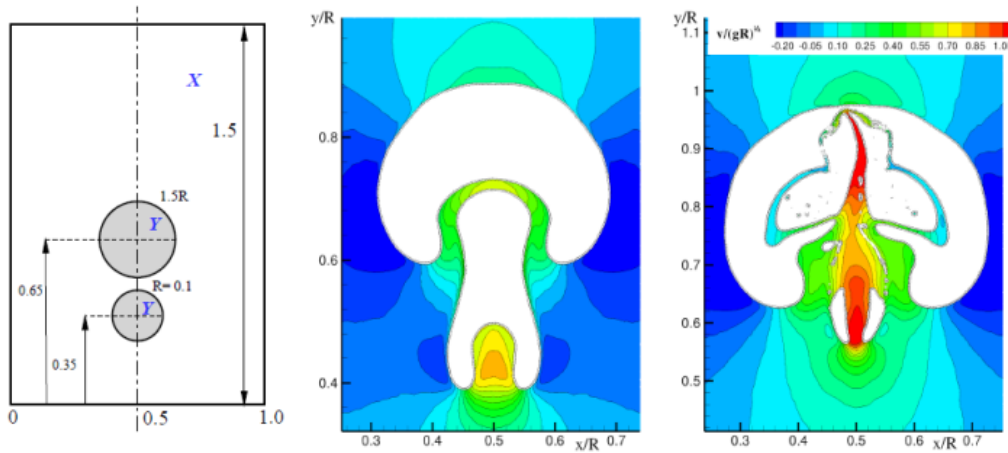
Fast aluminium beam impact



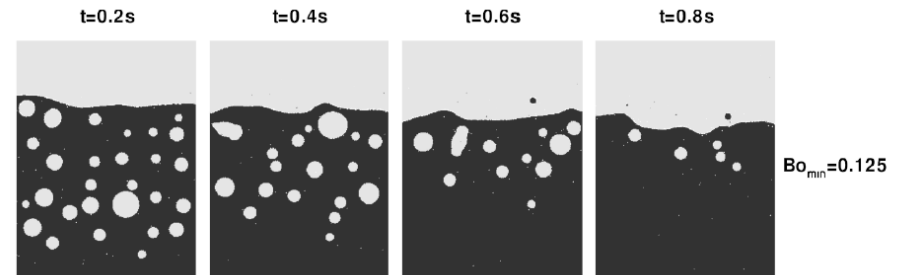
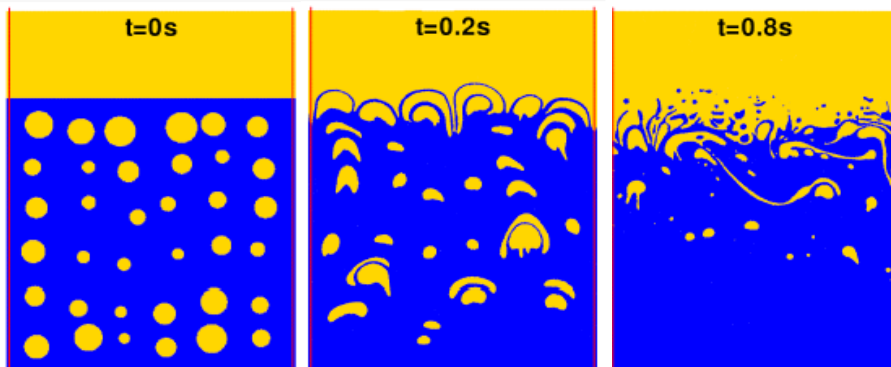
LNG membrane impact

Exemples d'application

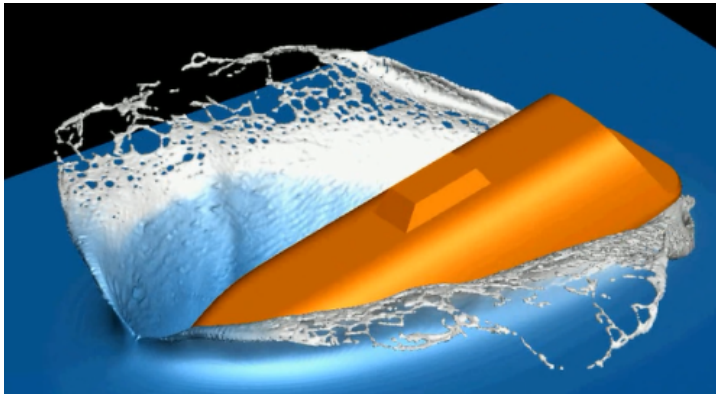
Complex bubbly flows (separation, atomisation)



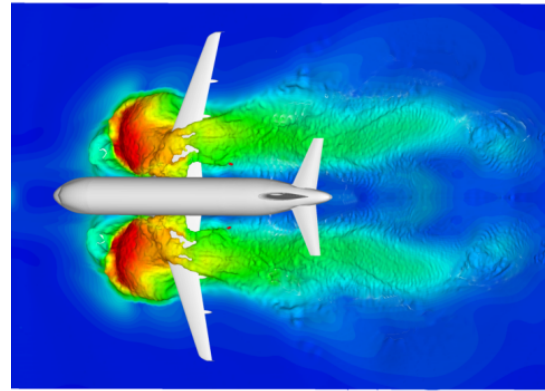
Water-oil separation (w/o or w/ surface tension)



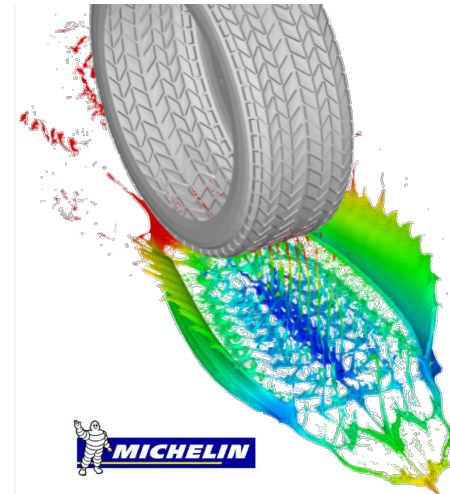
13. Industrial applications?



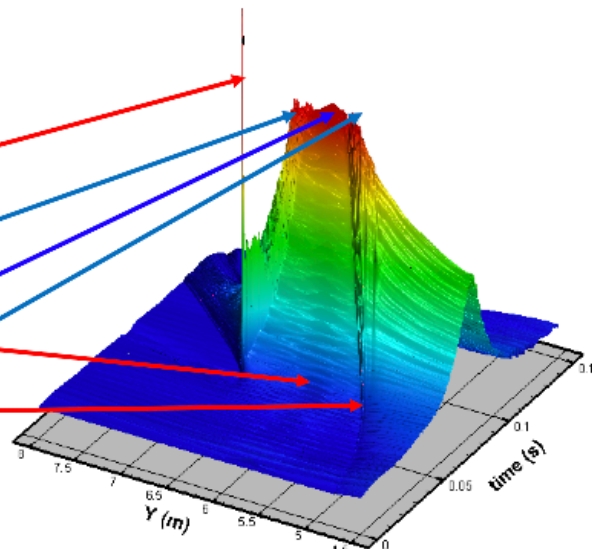
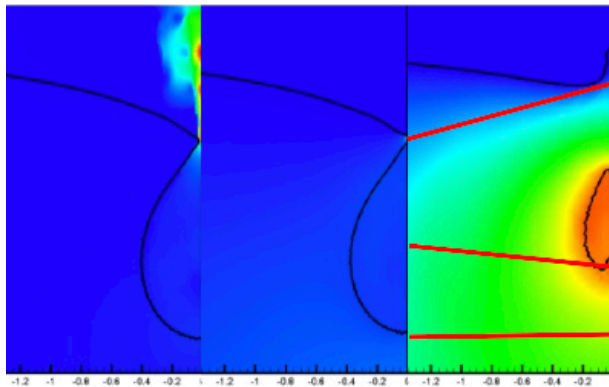
Lifeboat launching



Aircraft ditching



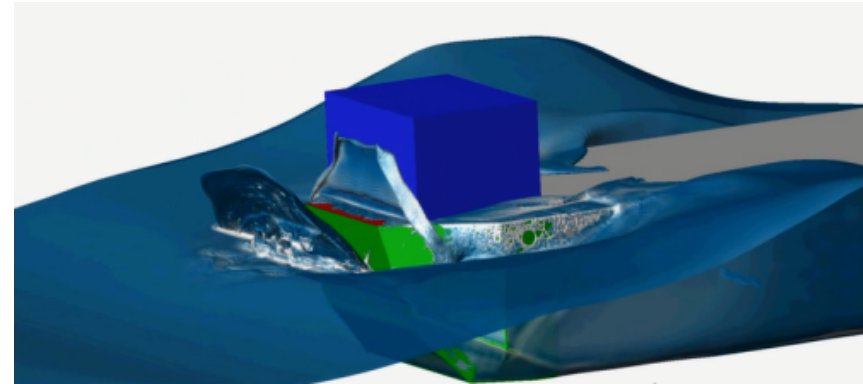
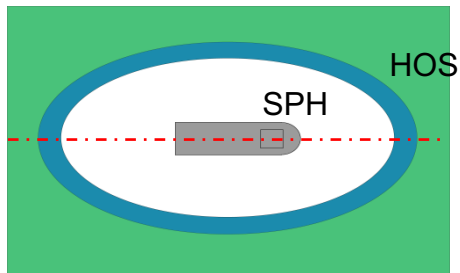
Hydroplaning



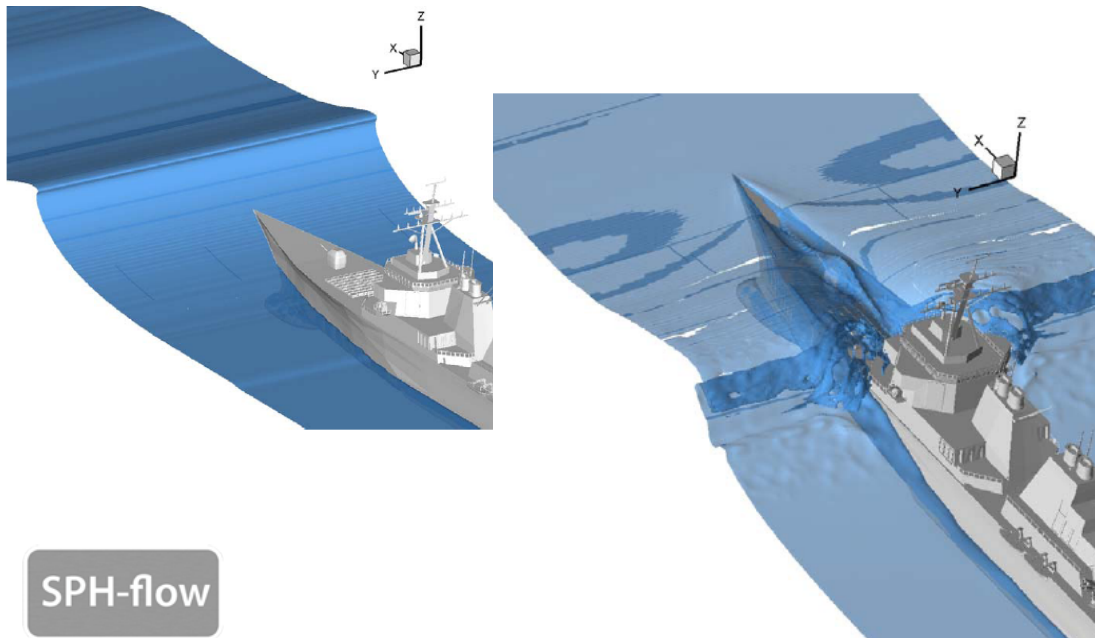
Liquid Natural Gas (LNG) sloshing

13. Industrial applications?

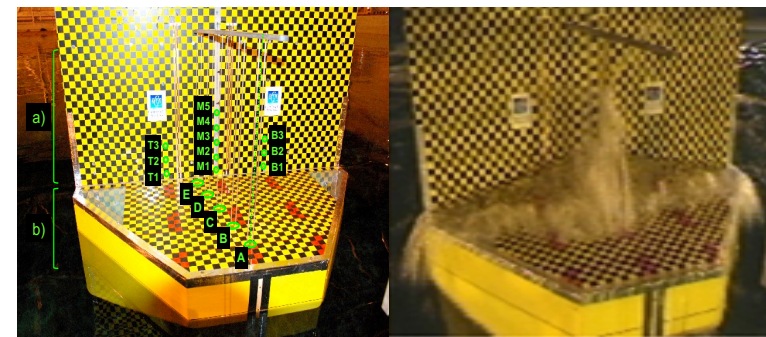
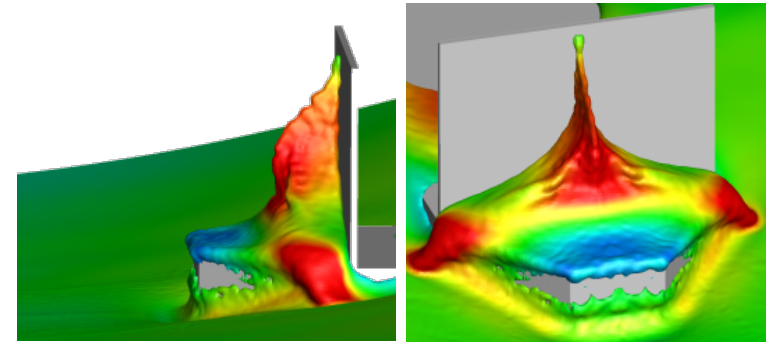
- FPSO in severe sea state coupling strategy:



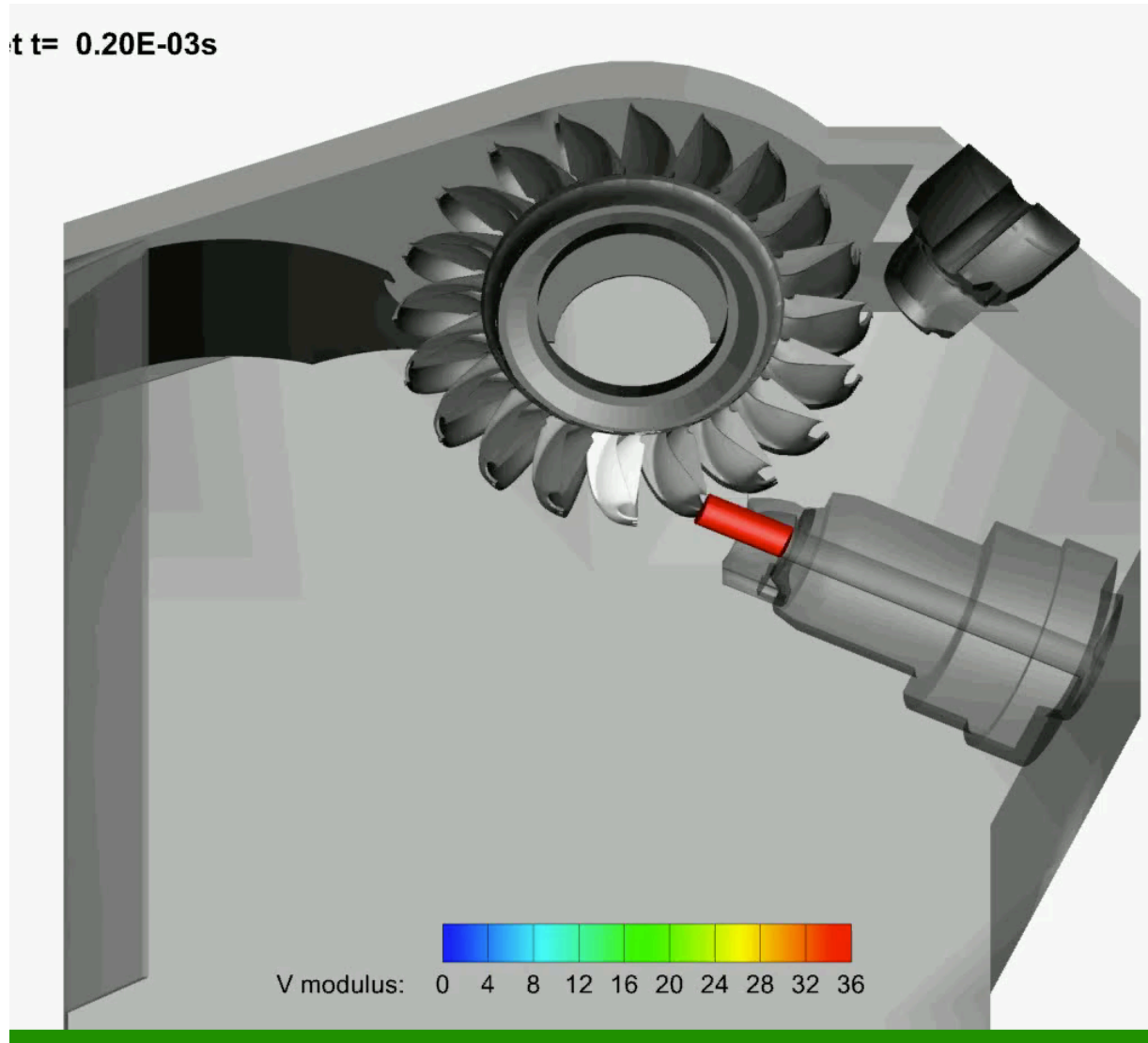
- Fregate facing a dimensioning wave



SPH-flow



13. Industrial applications?



SPH-flow

Pelton turbine

Conclusions



- A method growing for **flows with complex/multiple interfaces/bodies**
- The method extends also towards more and more **multiphysics fields**: ease to add PDEs in the system to solve and to have separated materials
- Numerical experience and understanding of the method fundamentals and numerical mechanisms is growing but a **difficulty remains in terms of numerical analysis**/applied maths on the method, slowing down progress towards higher-order, etc.
- A still **costly** method

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C&F incentive

C&F promotes benchmarking papers (cf. C&F website), so don't hesitate to go in this direction